

Charge Collection and Capacitance-Voltage analysis in irradiated n-type magnetic Czochralski silicon detectors

M. K. Petterson, H.F.-W. Sadrozinski, C. Betancourt
SCIPP UC Santa Cruz, 1156 High Street, 95064 CA, USA

M. Bruzzi, M. Scaringella, C. Tosi, A. Macchiolo
INFN and Università di Firenze, Italy

N. Manna, D. Creanza
Università di Bari, Italy

M. Boscardin, C. Piemonte, N. Zorzi
ITC, IRST, Povo, Trento, Italy

L. Borrello, A. Messineo
INFN Pisa, Italy

G. F. Dalla Betta
Università di Trento

Abstract

The depletion depth of irradiated n-type silicon microstrip detectors can be inferred from both the reciprocal capacitance and from the amount of collected charge. Capacitance-voltage (C-V) measurements at different frequencies and temperatures are being compared with the bias voltage dependence of the charge collection on an irradiated n-type magnetic Czochralski silicon detector. Good agreement between the reciprocal capacitance and the median collected charge is found when the frequency of the C-V measurement is selected such that it scales with the temperature dependence of the leakage current. Measuring C-V characteristics at prescribed combinations of temperature and frequency allows then a realistic estimate of the depletion characteristics of irradiated silicon strip detectors based on C-V data alone.

1. Introduction

The main operational parameter of a silicon microstrip detector is the bias voltage dependence of the charge collected at the electrodes when a particle is traversing the device. For a minimum ionizing particle (*mip*), the number of electron/hole pairs is generated uniformly in the detector with $73 \text{ e}/\mu\text{m}$ [1], and thus the most probable charge generated is linearly proportional to the thickness of the

depleted region. For this reason, the profile of the collected charge for *mips*, evaluated as a function of the applied reverse voltage ($CC(V)$), should exhibit the same trend as the reciprocal of the capacitance vs voltage ($C^{-1}(V)$), since the capacitance of a parallel-plate capacitor varies inversely with its thickness. (This is strictly true only as long as trapping effects can be neglected, which become very important at fluences above 10^{15} neq/cm² [2]).

Predicting the voltage dependence of the charge collection of irradiated sensors from straightforward electrical measurements facilitates the planning of experiments, and for many years C-V measurements have been used for that purpose, as the measurement of the $CC(V)$ is quite time consuming and experimentally more difficult with respect to a $C^{-1}(V)$ analysis. However in order to be able to quantify this comparison, it is very important to determine the experimental conditions for which the charge collection $CC(V)$ can be derived from the capacitance-voltage (C-V) characteristics. This is true in particular after fast hadron irradiation at high fluence, when detectors need to be operated cold to sufficiently decrease the leakage current I generated during irradiation:

$$\frac{I}{Volume} = \alpha \cdot \Phi \quad (1)$$

where $\alpha = 4 \cdot 10^{-17}$ A/cm for 1MeV neutrons (measured at 20°C after canonical annealing) [3]. The decrease in leakage current is required to permit higher operating voltage, decrease the noise and avoid self-heating effects. The current dependence on the absolute temperature T is:

$$I(T) \propto T^2 \exp\left(-\frac{E}{KT}\right) \quad (2)$$

The leakage current is mainly originated by the presence of generation-recombination deep levels close to mid-gap. Experimentally we find that the activation energy E is close to 0.6 eV for irradiated detectors [4]. It is well known that, after irradiation, radiation-induced deep traps are responsible for a strong dependence of the capacitance on the frequency and temperature. This work presents a simple and efficient method to quantitatively compare C-V at different temperatures with $CC(V)$ analyses at low temperature.

In section 2 the experimental set-up is described, and the theoretical framework of the correlation between temperature and frequency in C-V measurements is explained in section 3. In section 4 measurements on an irradiated p-on-n Magnetic Czochralski (MCz) Si detector are used to compare the charge collection with the inverse capacitances at different frequencies and temperatures.

2. Experimental Set-up

Measurements have been performed with microstrip detectors produced at IRST, Trento within the INFN SMART project [5]. They are processed on 300 μm thick n-type MCz Silicon wafers of about 1k Ωcm resistivity, with 4.5cm length and 50 μm pitch. The irradiation was performed at the Forschungszentrum Karlsruhe Cyclotron with 26 MeV protons up to 1 MeV neutron equivalent fluences in the range 10^{14} - 10^{15} cm^{-2} .

The charge collection measurements on microstrip detectors have been performed at SCIPP, UC Santa Cruz. The CC(V) investigations have been carried out with a ^{90}Sr source, using a small scintillator to trigger the binary data acquisition system (DAQ) [6] with a 200 ns shaping time. The C-V experimental set-up is based on a HP 4284A LCR meter coupled with the HP 16065A test fixture (modified to permit biasing with voltages up to 1000 V). Voltage sourcing and current monitoring are performed through a computer controlled Keithley 2410. In order to reduce the leakage current noise, CC(V) measurements have been carried out close to -10 $^{\circ}\text{C}$. Low temperature is obtained by spilling liquid nitrogen at the bottom of a thermally insulated box containing the sample holder. The temperature is sensed with a Pt resistor (Pt100) to a fraction of 1 $^{\circ}\text{C}$. Extraction of analogue information from the binary electronic read-out is accomplished by performing efficiency vs. threshold curves, which give the integral of the pulse-height spectrum. The following parameters are determined: the efficiency at a fixed threshold of 1 fC, used for example in ATLAS SCT, the median of the pulse height spectrum, i.e. the 50% efficiency point, and the most probable pulse height in the pulse height spectrum, obtained from the threshold curve by differentiation. The method is illustrated in Fig. 1 for an irradiated n-type MCz Si microstrip detector. More details on the experimental system are given in [6].

3. Frequency dependence of the capacitance in irradiated detectors

After heavy irradiation the silicon detector is highly compensated by the radiation-induced traps, and the shallow dopants (mainly P or B) initially present in the material are mostly removed or de-activated [3]. In this situation, the space charge of the device is essentially due only to the deep traps formed by irradiation. The emission coefficient e_n of these radiation induced traps is thermally activated through a Boltzmann factor depending on the distance in energy between their energy level and the edge of the conduction or valence band, E_t :

$$e_n \propto T^2 \exp\left(-\frac{E_t}{KT}\right) \quad (3)$$

During the C-V measurement an a.c. test signal is superimposed on the constant applied bias. Due to this a.c. signal, those deep traps that are close to the edge of the depletion region continuously switch between the depleted and the neutral regions. In the depleted region they are supposed to be empty and so contribute to the space charge of the device, while in the neutral region they are filled with charged carriers and thus neutral. If the frequency of the test signal is sufficiently low, i.e. comparable with the emission constant of the trap, this process of filling and re-emptying the traps will be measured by the device. Conversely, at sufficiently high test signal frequencies the measurement will be insensitive to the presence of the deep traps, so the material will appear essentially intrinsic. As a result, a flattening of the capacitance-voltage measurement will be observed.

As the emission coefficient is thermally activated (see eq.3) the test signal frequency f where this flattening is occurring depends on the temperature through: $f \sim e_n$ [7]. Based on this principle, the proper frequency $f(T)$ to be used at the temperature T can be scaled from the room temperature frequency $f(RT)$ using the respective emission coefficients e_n :

$$f(T) = \frac{e_n(T)}{e_n(RT)} f(RT) \quad (4)$$

Since in irradiated silicon, the traps with energy levels close to midgap are the main contributors to the radiation damage of the detector, the value of the activation energy in eq. (3) is approximately the same as in eq. (2), and the ratio of emission coefficients at two different temperatures in eq. (4) simply becomes the ratio of leakage currents measured at those two temperatures. Thus eq. (4) becomes:

$$f(T) \approx \frac{I(T)}{I(RT)} f(RT) \quad (5)$$

allowing to scale the frequency correctly when a change in the operational temperature is applied.

As a caveat, it is important to note that even before irradiation, C-V measurements on microstrip detectors exhibit a frequency dependence of the capacitance [8]. An example of frequency dependence of the capacitance observed before irradiation is shown in Fig. 2 for a microstrip detector, manufactured with n-type MCz Si. The measurements employ the usual algorithm of the LCR meter which represents the microstrip detector simply as one capacitance and one resistance in series. A wide plateau is observed in the range 0.5-100 kHz. In the figure the behavior of the capacitance as a function

of the frequency has been simulated by adding a term inversely proportional to the frequency, with a constant which scales with the observed resistance. A different value for the resistances of implants and the bulk will modify the measured frequency dependence of the capacitance. This effect is not due to the presence of radiation-induced defects in the material, but it is an artefact of the measurement procedure performed on the detector through the LCR meter, and can be corrected for by treating the strip detector correctly as a network of coupled distributed resistors and capacitors [8]. When occurring simultaneously, the two effects should be properly disentangled before proceeding to a correct comparison between the $CC(V)$ and $C^{-1}(V)$ profiles.

Our study will be focussed on the plateau range of frequency, where only the effect of deep traps is relevant. As a recent standardization of experimental measurement within the High Energy Physics Community has fixed the test signal frequency of a C-V characteristics to 10 kHz [9], we will consider this as the reference parameter at room temperature. The $C^{-1}(V)$ profile of a heavily irradiated detector measured at 10 kHz will appear considerably flattened going from room temperature to a lower temperature such as $-10\text{ }^{\circ}\text{C}$, as an effect of the frozen out of the carriers in the deep traps. As an example, Fig. 3 shows the $C^{-1}(V)$ for a proton irradiated n-type MCz Si detector ($1.7 \times 10^{14}\text{ cm}^{-2}$ 1MeV neutron equivalent): the shape of the curves measured at RT (22°C) and low temperature ($-10\text{ }^{\circ}\text{C}$) are indeed quite different.

The average ratio of the leakage current of this detector measured at different voltages and at two different temperatures ($T_R = 22.4^{\circ}\text{C}$ and $T_L = -11^{\circ}\text{C}$) is $\frac{I(T_R)}{I(T_L)} = 20.3 \pm 2.0$ (further details can be found in [10]). This scaling factor corresponds to an activation energy $E = 0.60 \pm 0.02\text{ eV}$ in eq. (3), in agreement with previous literature [4]. Table 1 lists the correct frequency to be used at low temperature T to best approximate a C-V characteristic measured at 10 kHz and room temperature (eq. (5)).

4. Comparison between charge collection and capacitance measurements in irradiated detectors

We are ultimately interested in comparing the $C^{-1}(V)$ with the charge collection $CC(V)$ measurements performed at low temperature. The $CC(V)$ curve obtained for a n-type MCz Si microstrip detector irradiated with 26 MeV protons up to the 1 MeV neutron equivalent fluence of $1.7 \times 10^{14}\text{ cm}^{-2}$ (after 5min annealing at 60°C) is shown in Fig.4. In the high voltage range the inefficiency of the detector is about 10%. We have superimposed the collected charge with the data of

three normalised $1/C$ vs. V measurements taken at the following pairs of frequency and temperature: 400 Hz and -11 °C; 10 kHz and 22 °C; 10 kHz and -11 °C. We observe that: (a) the 10 kHz curve at -11 °C is much too flat to describe the charge collection profile over any voltage range; (b) the 10 kHz curve at $+22$ °C is in agreement with the charge collection profile up to approximately 200 V, (for higher voltages the leakage current exceeds the limit of the C-V measurement system, and thus cannot be measured and compared with the $CC(V)$); (c) the low frequency $C^{-1}(V)$ curve measured at -11 °C and 400 Hz is fitting the charge collection profile much better over the entire voltage range.

An important parameter usually extracted from the C-V characteristics is the full depletion voltage V_{fd} , corresponding to the bias required to maximise the charge collected in the detector. In a detector characterised by a uniform doping, such that the effective doping concentration N_{eff} is constant, the $1/C^2$ vs. Voltage curve is a straight line for $V < V_{fd}$ and then is a constant for $V > V_{fd}$. In irradiated detectors, V_{fd} is usually determined as the intersection point of the two linear curves in the $1/C^2$ vs. voltage profile best fitting respectively the low range, where the sample is under-depleted, and the high voltage where over-depletion occurs. The $C^{-2}(V)$ curve for the irradiated microstrip detector is shown in Fig. 5 for two experimental combinations of the parameters (f,T): (10 kHz; 22 °C) and (400 Hz; -11 °C). Using the standard procedure for evaluating V_{fd} one obtains 48V and 140V at high and low temperature respectively, quite different. Using Fig. 4, we find first that these values corresponds to about the 45% and 80% of the maximum collected charge, while the maximum CC is actually achieved for $V \approx 350V$. This indicates that this standard method gives flawed results for the actual full depletion voltage at any frequency, but certainly at 10 kHz. Moreover, when performing $1/C$ measurements at the correct combination of frequency and temperature, the fraction of collected charge at any bias can be predicted.

The strong non-linearity of the $C^{-2}(V)$ at low T could be caused by a non-uniform doping of the device after irradiation. Although this analysis cannot reveal the nature of the radiation-induced deep traps responsible of this strong doping non-uniformity, we suppose that this effect is related to the occurring of a double junction in the irradiated detector [11]. Further analysis of this effect is in progress.

Conclusions

We have analysed the C-V curves of a n-type Si detector irradiated with 26MeV protons up to 1.7×10^{14} neq/cm² at different temperatures and frequencies, and have determined the correct scaling of the test signal frequency for capacitance measurements at low temperature to best reproduce the voltage dependence of the charge collection measurement. The best agreement in the entire voltage range is achieved when the C-V measurement is done at low temperature and low frequency (“LTLF”). The value of the required frequency can be determined by a proper scaling through the temperature dependence of the leakage current of the device. The $C^{-1}(V)$ profile done at LTLF can be used to evaluate the voltage required to maximise the collected charge. From the analysis of the median charge determined through an efficiency vs. threshold plot this value is achieved for $V \approx 350V$.

References

- [1] Yao WM et al., Review of Particle Physics, Chapter 27 Passage of Particles through Matter, J. Phys. G: Nucl. Part. Phys., 33 (2006) 265-267.
- [2] G. Casse et al., Nucl. Inst. Meth. A 518 (2004) 340-342.
- [3] M. Bruzzi, IEEE Trans. Nucl. Sci., 48, 4 (2001).
- [4] E. Barberis et al., Nucl. Inst. Meth. A 326 (1993) 373-380.
- [5] M. Bruzzi et al., Nucl. Inst. Meth. A 552 (2005) 20-26.
- [6] H. F.-W. Sadrozinski et al., IEEE Trans. Nucl. Sci., 51, 5, (2004) 2032.
- [7] E. Borchini, M. Bruzzi et al, Solid State Electronics, 42, 11, 2093 (1998)
- [8] E. Barberis et al , IEEE Trans. Nucl. Sci. 41, 1994, 785.
- [9] A. Chilingarov, Standardization of CV measurements, RD50 Workshop, 2004.
- [10] M. Bruzzi, Capacitance-Voltage analysis at different temperatures in heavily irradiated MCz silicon detectors, SCIPP 06/12 Report, 5/09/06. Available at: <http://www.scipp.ucsc.edu>
- [11] N. Manna et al., Results from the SMART INFN project, These Proceedings.

Figure and Table Captions

Fig. 1 (a) Efficiency vs. threshold charge for a n-type MCz Si microstrip detector irradiated with 26 MeV protons up to the fluence of $1.7 \times 10^{14} \text{ cm}^{-2}$ (5 min annealing at 60 °C) at 80 V and -10 °C. (b) Pulse height spectrum obtained by differentiation.

Fig. 2 Capacitance (a) and resistance (b) measured at room temperature in C_sR_s mode at 300 V for a p-on-n microstrip detector before irradiation.

Fig. 3 Reciprocal of the capacitance vs. voltage measured at 22°C and -10°C on a p-on-n irradiated MCz Si microstrip detector with a 10 kHz test signal frequency. The fluence of the irradiation was $1.7 \times 10^{14} \text{ cm}^{-2}$ (1 MeV n_{eq}), with no annealing performed after irradiation.

Fig. 4 Median value of the collected charge in an irradiated n-type MCz Si microstrip detector W187-S2 (Fluence of irradiation $1.7 \times 10^{14} \text{ cm}^{-2}$ 1 MeV n_{eq}), after 5 min annealing at 60 °C, compared to the normalised reciprocal capacitance measured at different temperatures with test signal of different frequencies.

Fig. 5 $1/C^2$ vs. voltage for the irradiated microstrip detector W187-S2 measured with different experimental parameters. Dashed lines are shown to demonstrate the standard method to determine the full depletion voltage.

Table 1: Frequency to be used at temperature T to best approximate the 10 kHz RT C-V characteristics