

## Evaluation of some integrals over solid angles—Part 2

In computing the total power liberated by an accelerating charge moving with velocity  $\vec{\beta} \equiv \vec{v}/c$  and acceleration  $\vec{a} \equiv d\vec{v}/dt$ , we need to compute three integrals over the solid angle  $\Omega$ ,

$$I_1 = \int \frac{(\hat{n} \cdot \vec{a})^2}{(1 - \vec{\beta} \cdot \hat{n})^5} d\Omega, \quad (1)$$

$$I_2 = \int \frac{d\Omega}{(1 - \vec{\beta} \cdot \hat{n})^3}, \quad (2)$$

$$I_3 = \int \frac{\hat{n} \cdot \vec{a}}{(1 - \vec{\beta} \cdot \hat{n})^4} d\Omega. \quad (3)$$

Let's begin with  $I_2$ . Choose the  $z$  axis to lie along the direction of  $\vec{\beta}$ . Then, it follows that  $\vec{\beta} \cdot \hat{n} = 1 - \beta \cos \theta$ , where  $\beta \equiv |\vec{\beta}|$ . Writing  $d\Omega = d \cos \theta d\phi$  and introducing  $w \equiv \cos \theta$ , it follows that

$$I_2 = 2\pi \int_{-1}^1 \frac{dw}{(1 - \beta w)^3} = \frac{2\pi}{\beta} \int_{1-\beta}^{1+\beta} \frac{dy}{y^3} = -\frac{\pi}{\beta} \left[ \frac{1}{(1 + \beta)^2} - \frac{1}{(1 - \beta)^2} \right], \quad (4)$$

after changing the integration variable to  $y = 1 - \beta w$ . Hence,

$$\boxed{I_2 = \frac{4\pi}{(1 - \beta^2)^2}} \quad (5)$$

Next, we can take the derivative of  $I_2$  with respect to  $\vec{\beta}$  by making use of eq. (2),

$$\frac{\partial I_2}{\partial \vec{\beta}} = 3 \int \frac{\hat{n} d\Omega}{(1 - \vec{\beta} \cdot \hat{n})^4}. \quad (6)$$

Thus, we can identify

$$I_3 = \frac{1}{3} \vec{a} \cdot \frac{\partial I_2}{\partial \vec{\beta}}. \quad (7)$$

We can evaluate the right hand side of eq. (7) by using the result obtained in eq. (4). Since eq. (4) is a function of  $\beta = |\vec{\beta}|$ , we can use the chain rule to write

$$\frac{\partial I_2}{\partial \vec{\beta}} = \frac{\partial \beta}{\partial \vec{\beta}} \frac{\partial I_2}{\partial \beta} = \frac{\vec{\beta}}{\beta} \frac{\partial I_2}{\partial \beta}. \quad (8)$$

To obtain the last step above, we noted that  $\beta = (\vec{\beta} \cdot \vec{\beta})^{1/2}$ . Hence, it follows that

$$\frac{\partial \beta}{\partial \vec{\beta}} = \frac{\partial}{\partial \vec{\beta}} (\vec{\beta} \cdot \vec{\beta})^{1/2} = \frac{1}{2} (\vec{\beta} \cdot \vec{\beta})^{-1/2} \frac{\partial}{\partial \vec{\beta}} (\vec{\beta} \cdot \vec{\beta}) = (\vec{\beta} \cdot \vec{\beta})^{-1/2} \vec{\beta} = \frac{\vec{\beta}}{\beta}. \quad (9)$$

Finally, we can use eq. (4) to evaluate  $\partial I_2 / \partial \beta$ ,

$$\frac{\partial I_2}{\partial \beta} = \frac{16\pi\beta}{(1 - \beta^2)^3}. \quad (10)$$

Hence, we end up with

$$\boxed{I_3 = \frac{16\pi}{3} \frac{\vec{a} \cdot \vec{\beta}}{(1 - \beta^2)^3}}. \quad (11)$$

Finally, we can use eq. (2) to obtain

$$\frac{\partial I_2}{\partial \beta_i} = 3 \int \frac{\hat{n}_i d\Omega}{(1 - \vec{\beta} \cdot \hat{n})^4}, \quad (12)$$

$$\frac{\partial^2 I_2}{\partial \beta_i \partial \beta_j} = 12 \int \frac{\hat{n}_i \hat{n}_j d\Omega}{(1 - \vec{\beta} \cdot \hat{n})^5}, \quad (13)$$

after two successive differentiations. Hence, we can identify,

$$I_1 = \frac{1}{12} \sum_{i,j} a_i a_j \frac{\partial^2 I_2}{\partial \beta_i \partial \beta_j}. \quad (14)$$

Note that eqs. (9) and (10) are equivalent to

$$\frac{dI_2}{d\beta_i} = \frac{16\pi\beta_i}{(1 - \beta^2)^3}. \quad (15)$$

The second derivative can now be easily evaluated with the help of eq. (9),

$$\begin{aligned} \frac{\partial^2 I_2}{\partial \beta_i \partial \beta_j} &= \frac{16\pi\delta_{ij}}{(1 - \beta^2)^3} + 16\pi\beta_i \frac{\beta_j}{\beta} \frac{\partial}{\partial \beta} \left( \frac{1}{(1 - \beta^2)^3} \right) \\ &= \frac{16\pi\delta_{ij}}{(1 - \beta^2)^3} + \frac{96\pi\beta_i\beta_j}{(1 - \beta^2)^4}. \end{aligned} \quad (16)$$

Hence, eq. (14) yields,

$$\boxed{I_1 = \frac{4\pi}{3} \frac{|\vec{a}|^2}{(1 - \beta^2)^3} + \frac{8\pi(\vec{a} \cdot \vec{\beta})^2}{(1 - \beta^2)^4}}. \quad (17)$$

There is an alternative technique for evaluating  $I_1$  and  $I_3$ . We can define the following two integrals,

$$J_{ij} = \int \frac{\hat{n}_i \hat{n}_j d\Omega}{(1 - \vec{\beta} \cdot \hat{\mathbf{n}})^5}, \quad (18)$$

$$K_i = \int \frac{\hat{n}_i d\Omega}{(1 - \vec{\beta} \cdot \hat{\mathbf{n}})^4}. \quad (19)$$

By the covariance properties of Euclidean tensors, it follows that

$$J_{ij} = c_1 \delta_{ij} + c_2 \beta_i \beta_j, \quad (20)$$

$$K_i = \kappa \beta_i. \quad (21)$$

Consider first the evaluation of  $K_i$ . Multiplying by  $\beta_i$  and summing over  $i$  yields

$$\kappa \beta^2 = \int \frac{\vec{\beta} \cdot \hat{\mathbf{n}} d\Omega}{(1 - \vec{\beta} \cdot \hat{\mathbf{n}})^4}. \quad (22)$$

The integral above is now easily evaluated by employing the same method used to obtain eq. (4). Thus, we can obtain an explicit expression for  $\kappa$ . I will leave it as an exercise for the reader to show that

$$\kappa = \frac{16\pi}{3} \frac{1}{(1 - \beta^2)^3}. \quad (23)$$

Likewise, to evaluate  $J_{ij}$ , we first multiply by  $\delta_{ij}$  and sum over  $i$  and  $j$  to get one equation. A second equation is obtained by multiplying by  $\beta_i \beta_j$  and summing over  $i$  and  $j$ . Thus, we get two equations for the two unknowns  $c_1$  and  $c_2$ ,

$$3c_1 + c_2 \beta^2 = \int \frac{d\Omega}{(1 - \vec{\beta} \cdot \hat{\mathbf{n}})^5}, \quad (24)$$

$$c_1 \beta^2 + c_2 \beta^4 = \int \frac{(\vec{\beta} \cdot \hat{\mathbf{n}})^2 d\Omega}{(1 - \vec{\beta} \cdot \hat{\mathbf{n}})^5}. \quad (25)$$

Again, the two integrals above are easily evaluated by employing the same method used to obtain eq. (4). One can now solve for  $c_1$  and  $c_2$ . I will leave it as an exercise for the reader to carry out the remaining computations to obtain,

$$c_1 = \frac{4\pi}{3} \frac{1}{(1 - \beta^2)^3}, \quad c_2 = \frac{8\pi}{(1 - \beta^2)^4}. \quad (26)$$

Finally, we obtain

$$I_1 = \sum_{i,j} a_i a_j J_{ij}, \quad I_3 = \sum_i a_i K_i. \quad (27)$$

Using eqs. (20), (21), (23) and (26), we recover the results obtained in eqs. (17) and (11), respectively.