

Introduction to SU(5) and SO(10) Grand Unification

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Overview

- 1 Motivation for Grand Unification
- 2 SM introduction (first generation)
- 3 SU(5) Unification
- 4 SO(10) Unification (briefly)

Motivation

- Can the large number of parameters inserted into the standard model be reconciled?
- Is it possible to unify SM forces, account for seemingly arbitrary coupling constants?
- Likewise can the fermions be united?
- Why are the hypercharge values of multiplets the way they are?
- Similarly, why are the electric charges what they are?
- SM incompatible with massive neutrinos

Isospin

- Analogous to spin- $\frac{1}{2}$, described by SU(2)
- Strong isospin states up/down equivalent to spin- $\frac{1}{2}$ projections: $\pm\frac{1}{2}$
- Quite descriptive due to the proximity of the masses of the u,d quarks
- Strong interaction acts on multiplets of the total isospin, e.g.
 $2 \otimes 2 = 3 \oplus 1: \pi^{0,\pm} \oplus \eta$

Hypercharge

- Utility in assigning these particle families a particular number: Y
- Gell-Mann-Nishijima formula: $Q = I_3 + \frac{Y}{2}$

Weak isospin

- Extend idea to leptons, create second doublet: $\begin{pmatrix} \nu_L \\ e_L^- \end{pmatrix}$
- Weak isospin mediated by W^\pm, Z^0
- Observed asymmetry in chirality, only left-handed particles transform non-trivially under $SU(2)$
- Extend hypercharge as well, Gell-Mann-Nishijima still holds if the weak isospin doublet has $Y = -1$
- Right-handed particles have no weak isospin projections

Gauge bosons, symmetry breaking

- SM gauge group: $SU(3) \otimes SU(2) \otimes U(1)$
- $SU(3)$ gauge bosons: 8 massless gluons
- Unified electroweak gauge bosons (massless): $W^{0,\pm}, B$, with γ, Z^0 being linear combinations of W^0 and B
- $SU(2) \otimes U(1)$ symmetry broken, electromagnetic and weak force behave differently, W^\pm, Z^0 bosons become massive

SM Summary

- 16 total fermions in $SU(3) \otimes SU(2) \otimes U(1)$
- $\nu_R = (1 \otimes 1)_0$
- $e_R = (1 \otimes 1)_{-2}$
- Left-handed lepton doublet $(1 \otimes 2)_{-1}$
- Right-handed down quark triplet $d_R = (3 \otimes 1)_{-\frac{2}{3}}$
- Right-handed up quark triplet $u_R = (3 \otimes 1)_{\frac{4}{3}}$
- Left-handed quark sextuplet $Q_L = (3 \otimes 2)_{\frac{1}{3}}$
- 16 fermions and 16 anti-fermions: $16 \oplus \bar{16} = 32$

Why SU(5)?

- Sensible to start with a rank 4 group
- SM has 4 commuting generators, $T_{3,8}$ from SU(3), σ_3 from SU(2) and the U(1) generator
- Is it possible to embed the SM into SU(5)?
- By construction, the diagonal generators can be made equivalent to those of the SM

Fitting $SU(3) \otimes SU(2) \otimes U(1)$ in $SU(5)$

- Represent matrices in $SU(5)$ as block diagonal, with block elements that are $SU(3)$, $SU(2)$ matrices multiplied by an appropriate $U(1)$ factor

- $$A = \left(\begin{array}{c|c} \alpha^{-2}M & 0 \\ \hline 0 & \alpha^3N \end{array} \right), \alpha \in U(1), M \in SU(3), N \in SU(2)$$

- kernel of $\phi : (\alpha, M, N) \mapsto A$ is $(\alpha, \alpha^2, \alpha^{-3})$
- Satisfying determinant condition requires that α be a sixth root of unity: \mathbb{Z}_6

Fitting $SU(3) \otimes SU(2) \otimes U(1)$ in $SU(5)$ cont.

- To be able to embed the SM group (G_{SM}) into $SU(5)$, the kernel must act trivially
- With some algebra, this yields constraints on the hypercharge
- e.g. for left-handed quarks, $Y = 2k + \frac{1}{3}$, $k \in \mathbb{Z}$
- Structure introduced to the hypercharges
- Via Gell-Mann-Nishijima, electric charge determined

Where the first generation particles fit

- Examining the structure of that matrix A , its action would be on a vector made up of a doublet and a triplet
- The fundamental $5 = (1 \otimes 2)_1 \oplus (3 \otimes 1)_{-\frac{2}{3}}$
- Utilise $\bar{5} = (1 \otimes 2)_{-1} \oplus (\bar{3} \otimes 1)_{\frac{2}{3}}$, identify with doublet of leptons: $(\nu, e)^T$ and triplet of \bar{d}
- The antisymmetric part of the product of two 5s: $5 \otimes_A 5 = 10$ can describe the remainder
- $10 = (1 \otimes 1)_2 \oplus (3 \otimes 2)_{\frac{1}{3}} \oplus (\bar{3} \otimes 1)_{-\frac{4}{3}}$
- Identify \bar{e} , quark sextuplet, \bar{u} triplet

Gauge bosons

- $5^2 - 1 = 24$ generators

-

$$T^a = \begin{pmatrix} \frac{\lambda^a}{2} & 0 \\ 0 & 0 \end{pmatrix}, T^i = \begin{pmatrix} 0 & 0 \\ 0 & \frac{\sigma^i}{2} \end{pmatrix}$$

- λ^a and σ^i are the Gell-Mann and Pauli matrices respectively.
- for $a = 3, 8$ and $i = 3$, there are 3 diagonal matrices. We expect one more
- $\tilde{Y} = \frac{1}{\sqrt{60}} \text{diag}(-2, -2, -2, 3, 3)$
- $24 = (8 \otimes 1)_0 \oplus (1 \otimes 3)_0 \oplus (1 \otimes 1)_0 \oplus (\bar{3} \otimes 2)_{\frac{5}{6}} \oplus (3 \otimes 2)_{-\frac{5}{6}}$
- Twelve off-diagonal matrices, linear combinations of which give 12 additional bosons
- These X,Y bosons can mediate processes that don't conserve baryon or lepton number

Proton Decay

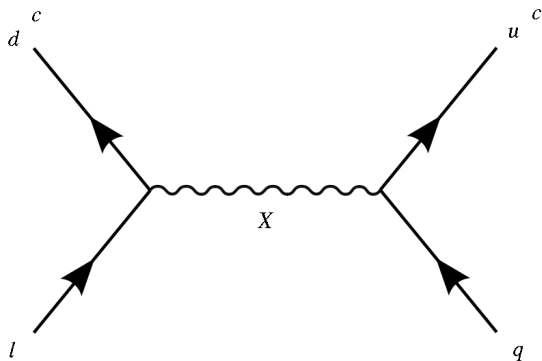


Figure: Retrieved from wikipedia. public domain.

Proton Decay

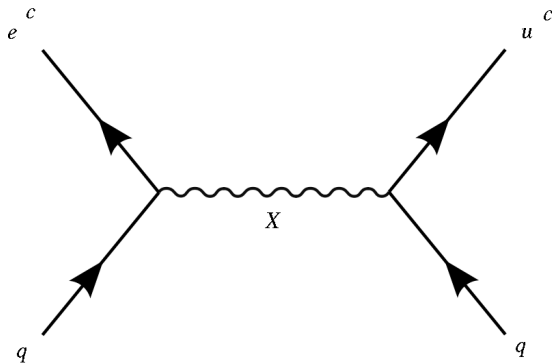
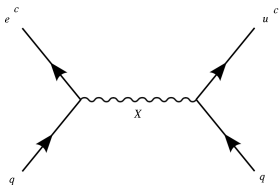


Figure: Retrieved from wikipedia. public domain.

Proton Decay



- $\Gamma(qqq \rightarrow q\bar{q}l = \pi^0 e^+) \approx \frac{m_p^5}{M_{GUT}^4}$
- Experimental searches for proton decay constrain $M_{GUT} \gtrsim 10^{16} \text{ GeV}$

Higgs, symmetry breaking, doublet-triplet problem

- For the Higgs, again use fundamental 5 containing a triplet and a doublet.
- Symmetry breaking $SU(5) \rightarrow SU(3) \otimes SU(2) \otimes U(1)$
- Doublet energy goes to the weak energy scale (100 GeV)
- Energy of triplet would remain at GUT energy scale, 10^{14} GeV
difference: fine tuning problem
- $SO(10)$ encounters same problem

Weinberg angle

- $\sin^2(\theta_W) = \frac{g'^2}{g^2+g'^2}$ where g' , g are the couplings for hypercharge and weak isospin respectively
- Denotes the rotation of the basis vectors γ, Z^0 in the W^0, B plane after electroweak symmetry breaking
- Can express it in terms of single SU(5) coupling g , replacing the g in the above equation.
- Inspect the action of \tilde{Y} on the states in the $\bar{5}$ and we can read off $\tilde{Y} = \frac{\sqrt{60}}{6} Y$
- Comparing Y to \tilde{Y} yields $(g/g')^2 = \frac{5}{3}$
- Predicts $\sin^2(\theta_W^{GUT}) = \frac{3}{8}$

SU(5) summary

- A single gauge coupling
- Structure of hypercharge, electric charge
- Energy scale large due to proton stability
- Weinberg angle prediction is motivating but simple scheme untenable
- Not all SM particles embedded into single irrep
- Fine tuning problem present

Pati-Salam model

- Idea: promote right handed singlets to doublets
- Weak isospin now acts like $SU(2) \otimes SU(2)$, with each $SU(2)$ acting on a doublet of different handedness
- Additionally, can think of lepton number as a fourth colour.
 $SU(3) \rightarrow SU(4)$
- Full model: $SU(4) \otimes SU(2) \otimes SU(2)$

- What can be gained by adding another commuting generator?
- B-L conservation is a naturally motivated choice: gives rise to Pati-Salam
- Can represent SM fermions inside a single irrep: $16 = 10 \oplus \bar{5} \oplus 1$
- All particles inside a single irrep

Breaking

- Can break symmetry in various ways, some examples:
- $SO(10) \rightarrow SU(3) \otimes SU(2) \otimes U(1)$
- $SO(10) \rightarrow SU(4) \otimes SU(2) \otimes SU(2)$
- $SO(10) \rightarrow SU(3) \otimes SU(2) \otimes SU(2) \otimes U(1)$
- $SO(10) \rightarrow SU(5) \otimes U(1)$

- We saw that the 5 of SU(5) was an SU(3) colour triplet and an SU(2) isospin doublet
- With $16 \oplus \bar{16} = 32$ SM fermions (and anti-fermions), can represent every particle in a 5 element vector, whose elements are binary flags for some kind of charge: $2^5 = 32$
- Can think of a binary flag as some SU(2) projection
- $\psi = (r, g, b, up, down)^T$

2^5 representation

| ψ | Y | r | g | b | u | d |
|---------|------|---|---|---|---|---|
| ν^c | 0 | 0 | 0 | 0 | 0 | 0 |
| e^c | 2 | 0 | 0 | 0 | 1 | 1 |
| u_r | 1/3 | 1 | 0 | 0 | 1 | 0 |
| d_r | 1/3 | 1 | 0 | 0 | 0 | 1 |
| u_g | 1/3 | 0 | 1 | 0 | 1 | 0 |
| d_g | 1/3 | 0 | 1 | 0 | 0 | 1 |
| u_b | 1/3 | 0 | 0 | 1 | 1 | 0 |
| d_b | 1/3 | 0 | 0 | 1 | 0 | 1 |
| u_r^c | -4/3 | 0 | 1 | 1 | 0 | 0 |
| u_g^c | -4/3 | 1 | 0 | 1 | 0 | 0 |
| u_b^c | -4/3 | 1 | 1 | 0 | 0 | 0 |
| d_r^c | 2/3 | 0 | 1 | 1 | 1 | 1 |
| d_g^c | 2/3 | 1 | 0 | 1 | 1 | 1 |
| d_b^c | 2/3 | 1 | 1 | 0 | 1 | 1 |
| ν | -1 | 1 | 1 | 1 | 0 | 1 |
| e | -1 | 1 | 1 | 1 | 1 | 1 |

Summary

- Various ways to unite leptons and quarks
- Hypercharge seen to be less arbitrary, electric charge follows
- Still a large number of parameters to be inserted
- SU(5): $M_{GUT} \gtrsim 10^{16} \text{ GeV}$
- SU(5) and SO(10) encounter fine tuning issues
- Extending these theories with SUSY possible, though thoroughly outside the scope here

References

- Johannes, B. (2016). Grand Unified Theories for Pedestrians. Presentation, DESY at Zeuthen.
- Malinsk, M. (2009). 35 years of GUTs - where do we stand?. Presentation, Royal Institute of Technology, Stockholm.
- Dine, M. (2019) Grand Unification and Monopoles [PowerPoint presentation].
- W.-M. Yao et al. (Particle Data Group), J. Phys. G33, 1 (2007)
- Georgi, H. (2018). Lie algebras in particle physics: from isospin to unified theories. CRC Press.