

Signatures of New Hidden Sectors at the LHC

Echoes of a hidden valley at hadron colliders.

M.J.S. & K. M. Zurek , hep-ph/0604261 , Phys.Lett.B651:374-379

Discovering the Higgs through highly-displaced vertices.

M.J.S. & K. M. Zurek , hep-ph/0605193

Possible effects of a hidden valley on supersymmetric phenomenology.

M.J.S., hep-ph/0607160

Phenomenology of hidden valleys at hadron colliders.

Han, Si, Zurek & M.J.S., arXiv/0712.2041

Why Unparticle models with mass gaps are examples of hidden valleys.

M.J.S., arXiv/0801.0629

Several papers in preparation...

See also Ciapetti, Lubatti, Dionisi...M.J.S. ATLAS note

Matthew Strassler
Rutgers University

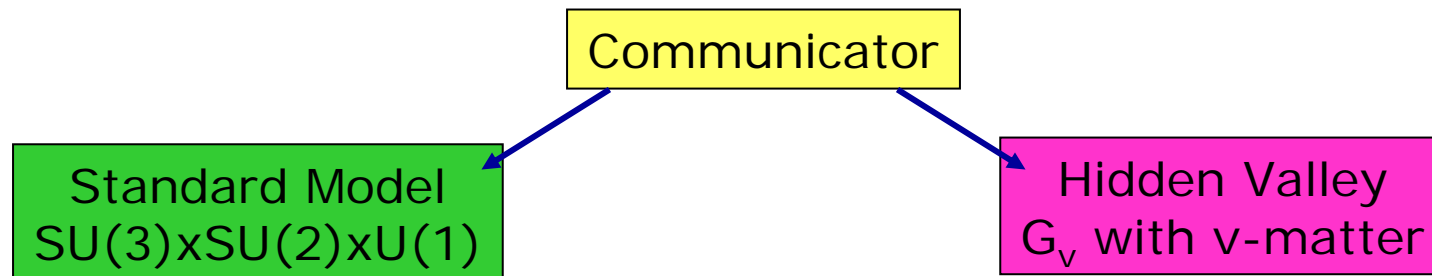
Plan of the Talk

- Basic structure and phenomenology of Hidden Valley scenario
- Relation to unparticle models with mass gaps
- Subtleties and phenomenology of such models
- If time, novel signals of hidden-sector strong dynamics

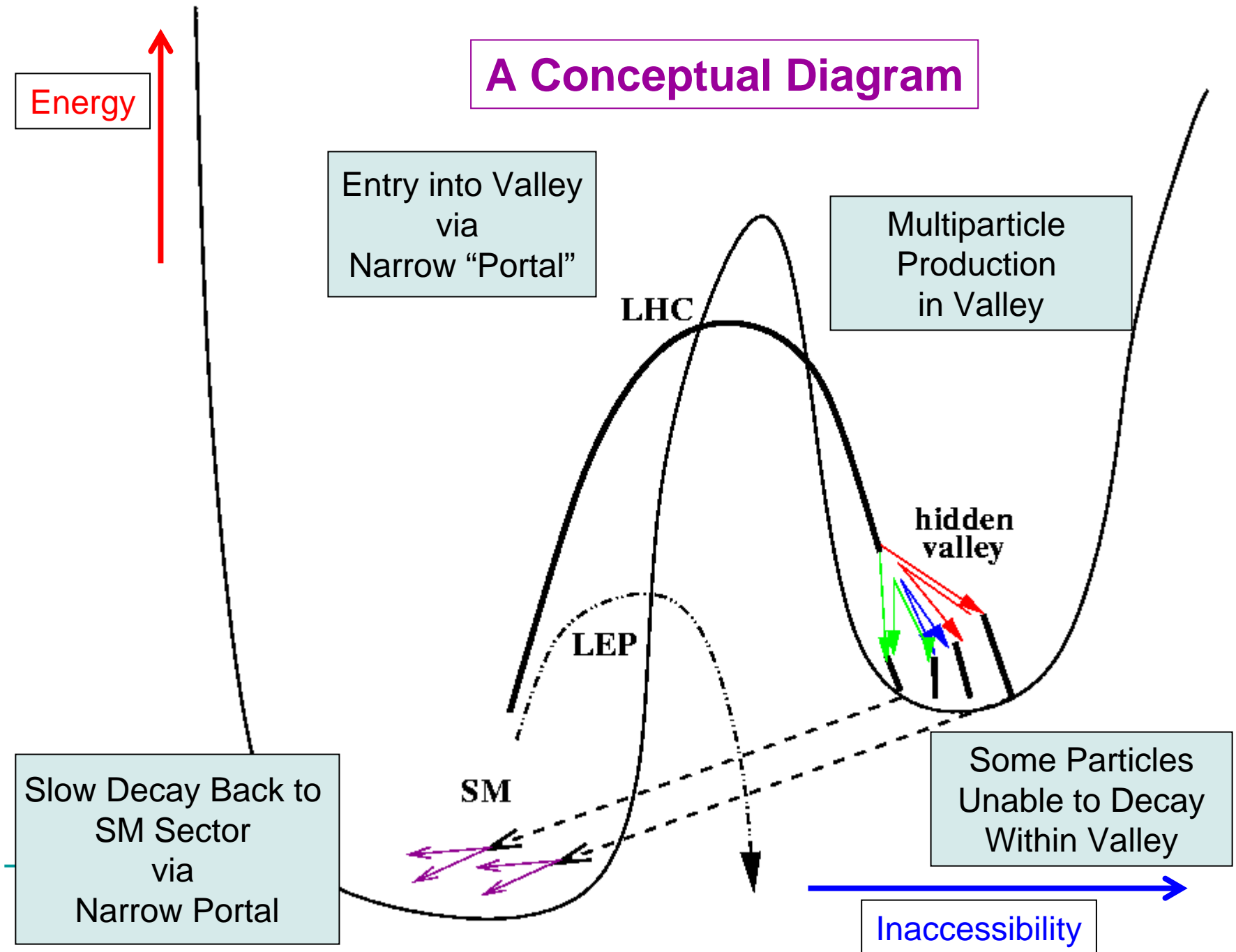
Hidden Valley Scenario (w/ K. Zurek)

hep-ph/0604261

- A scenario: **Large Meta-Class of Models**
- Motivation: in a moment
- Basic structure:

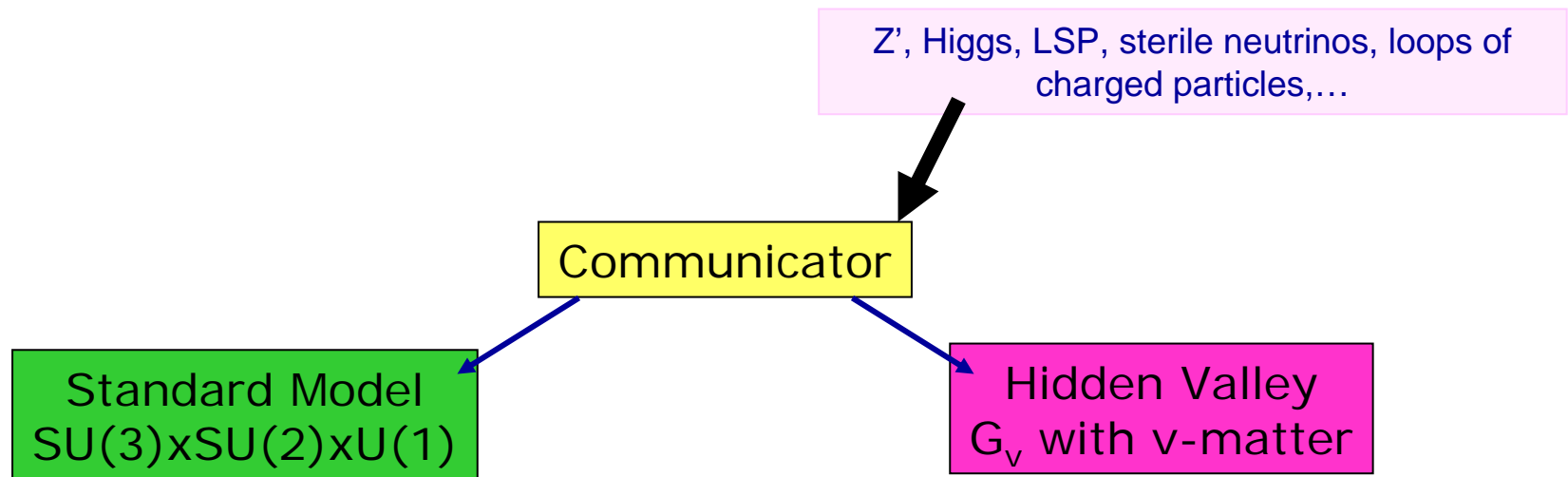


A Conceptual Diagram



Hidden Valley Models (w/ K. Zurek)

hep-ph/0604261



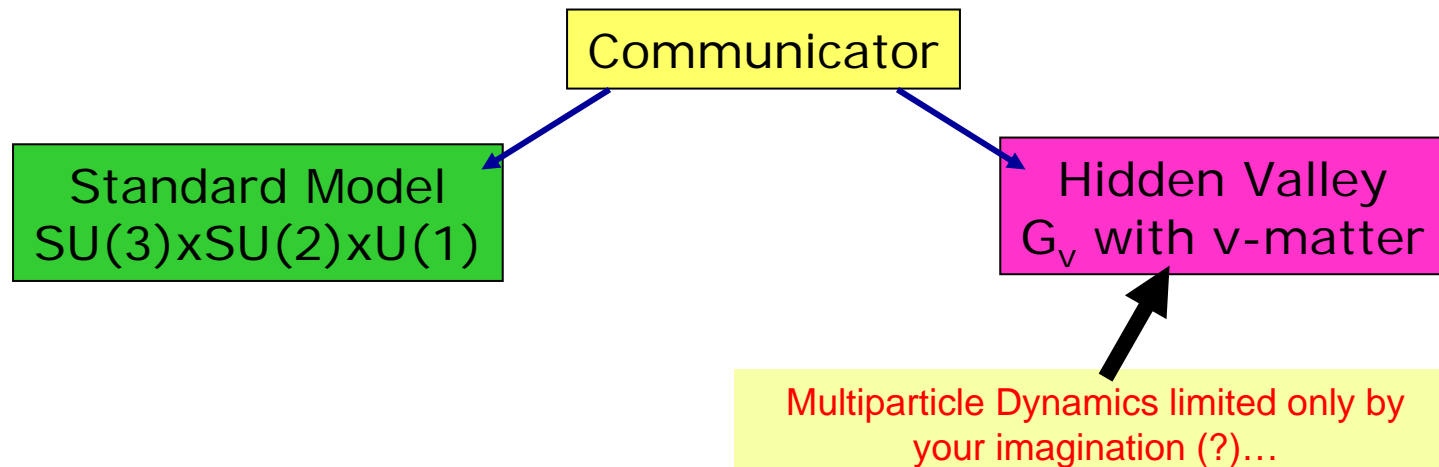
Hidden Valley Models (w/ K. Zurek)

hep-ph/0604261

Vast array of possible v-sectors...

QCD-like theory : F flavors and N colors
QCD-like theory : only heavy quarks
QCD-like theory : adjoint quarks
Walking-Technicolor-like theory
Pure-gluon theory
...

N=4 SUSY \rightarrow N=1 (N=1*)
RS or KS throat
Almost-supersymmetric N=1 model
Moose/Quiver model
Broken/Tumbling SU(N) theory
...



Motivation and Approach

■ Why the Hidden Valley Scenario?

- ❑ Extra sectors common in string theory, SUSY breaking, Extra dims, etc.
- ❑ Incredibly exciting if found: new particles, forces, dynamics [possibly strong]
- ❑ Can drastically change phenomenology of SUSY/Extra Dims/etc.
- ❑ Dark matter, early universe cosmology, astrophysics ?

■ The challenge of the Hidden Valley Scenario

- ❑ Weak experimental constraints!
- ❑ Vast array of possibilities
- ❑ **Phenomenology very challenging for hadron colliders – urgent!!**

■ Our approach:

- ❑ Find **characteristic predictions** of large classes of models at once
 - ❑ Produce search strategies, Monte Carlo tools that experimentalists can use now
-

Common Predictions of HV Scenario

MJS + Zurek hep-ph/0604261

- “Common predictions of Supersymmetry are missing transverse momentum, high p_T jets, medium p_T leptons”
 - But of course not all SUSY models do this
 - And some other non-SUSY models also do this
- Common predictions of Hidden Valley models are...

Common Predictions of HV Scenario

MJS + Zurek hep-ph/0604261

■ New Metastable Neutral Excitations “X”

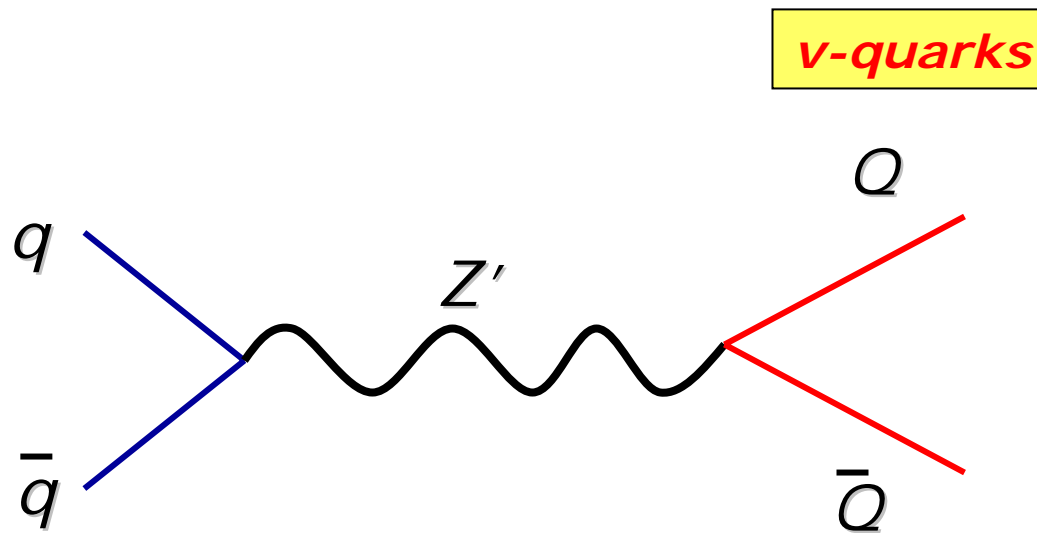
- Typically **more than one** X in v-sector
- Possibly quite **light** (mass $< M_Z$)
- Often long-lived; **displaced vertices, missing transverse energy**

■ Moderate- to High-Multiplicity Final States

- Exceptionally **busy final states** possible
- Unusual **event shapes**
- Jets \Leftrightarrow partons **breaks down**
- Very **large fluctuations** in appearance of events

A Confining Hidden Valley

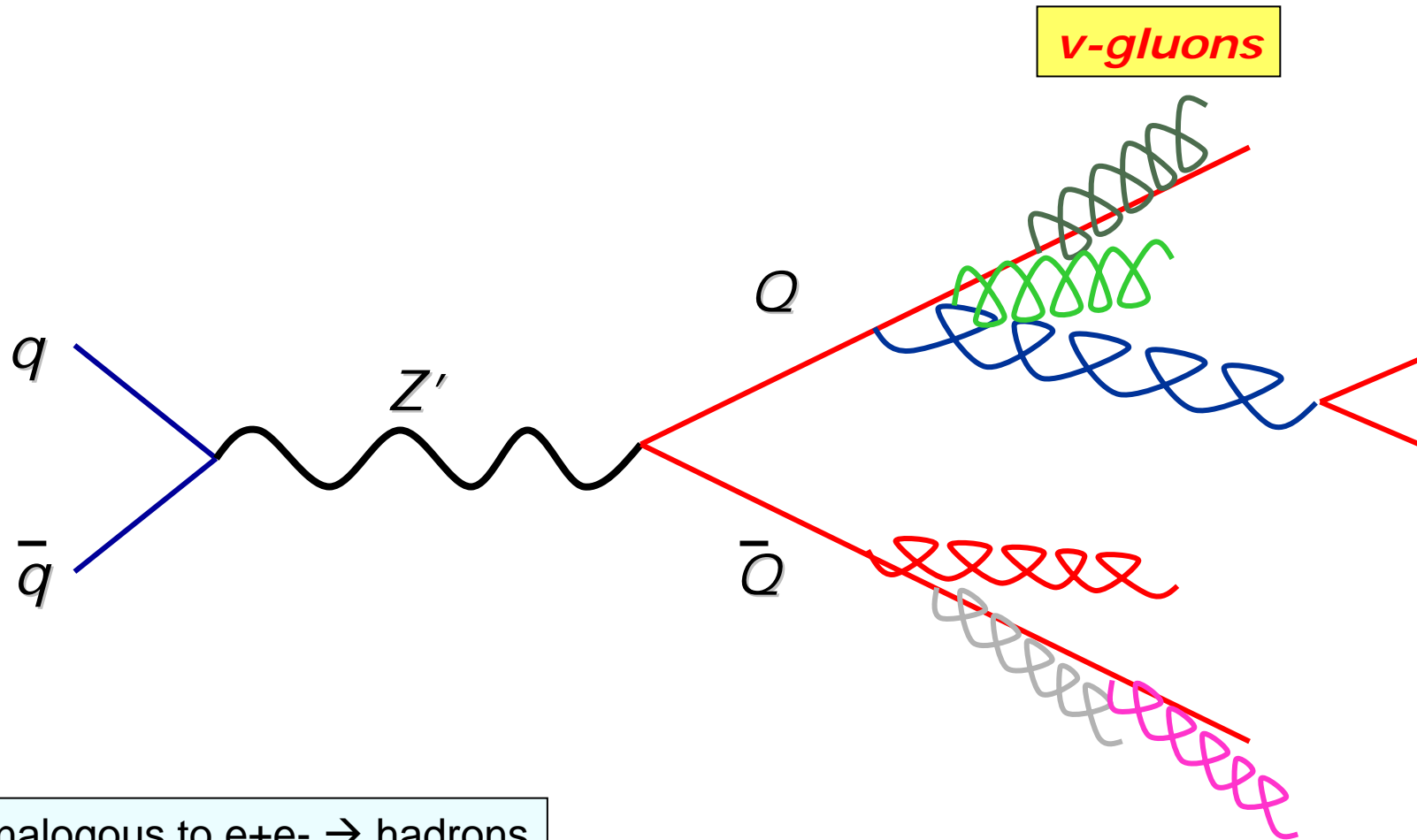
MJS + Zurek hep-ph/0604261



Analogous to $e^+e^- \rightarrow \text{hadrons}$

A Confining Hidden Valley

MJS + Zurek hep-ph/0604261

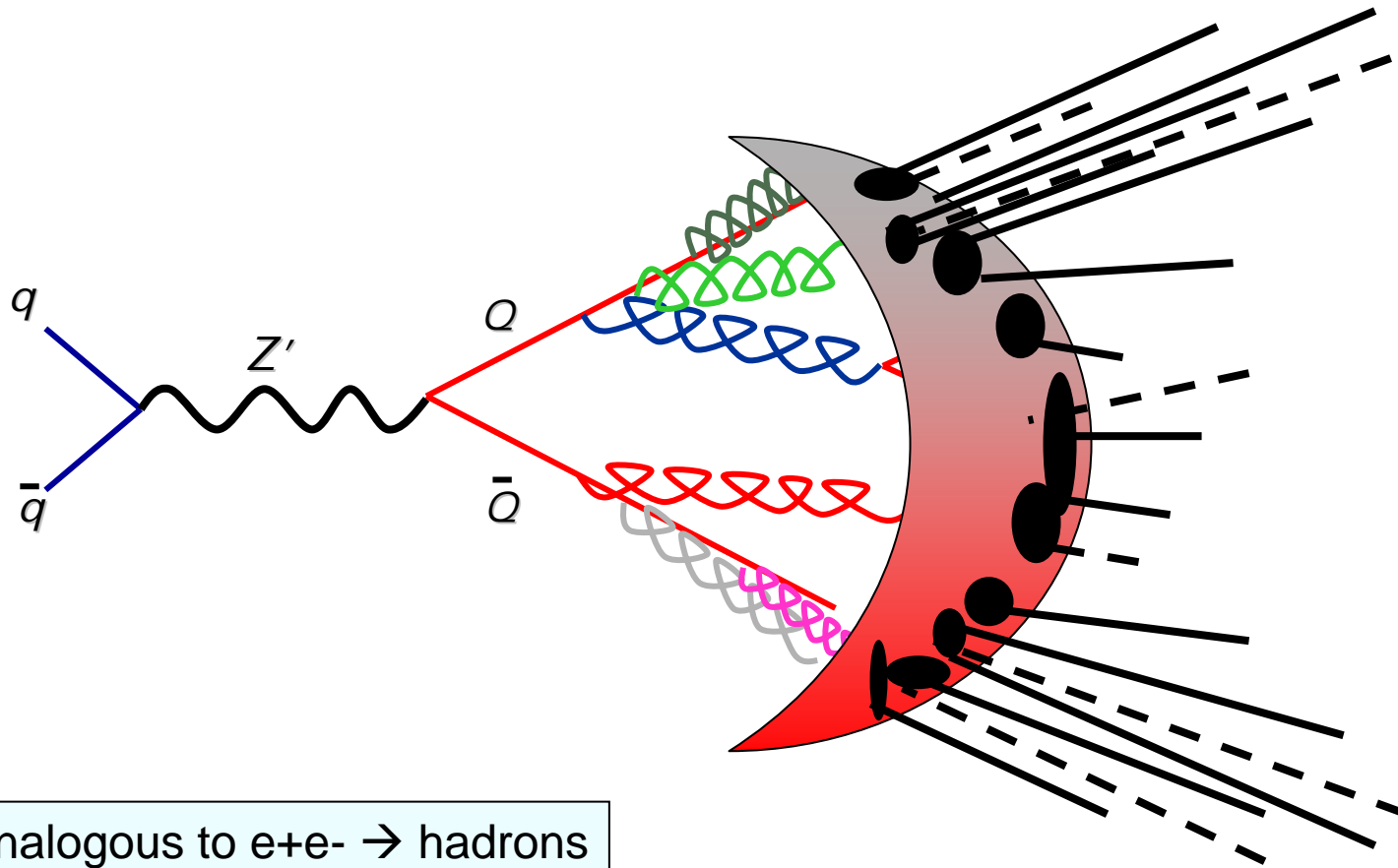


Analogous to $e^+e^- \rightarrow \text{hadrons}$

A Confining Hidden Valley

MJS + Zurek hep-ph/0604261

v-hadrons

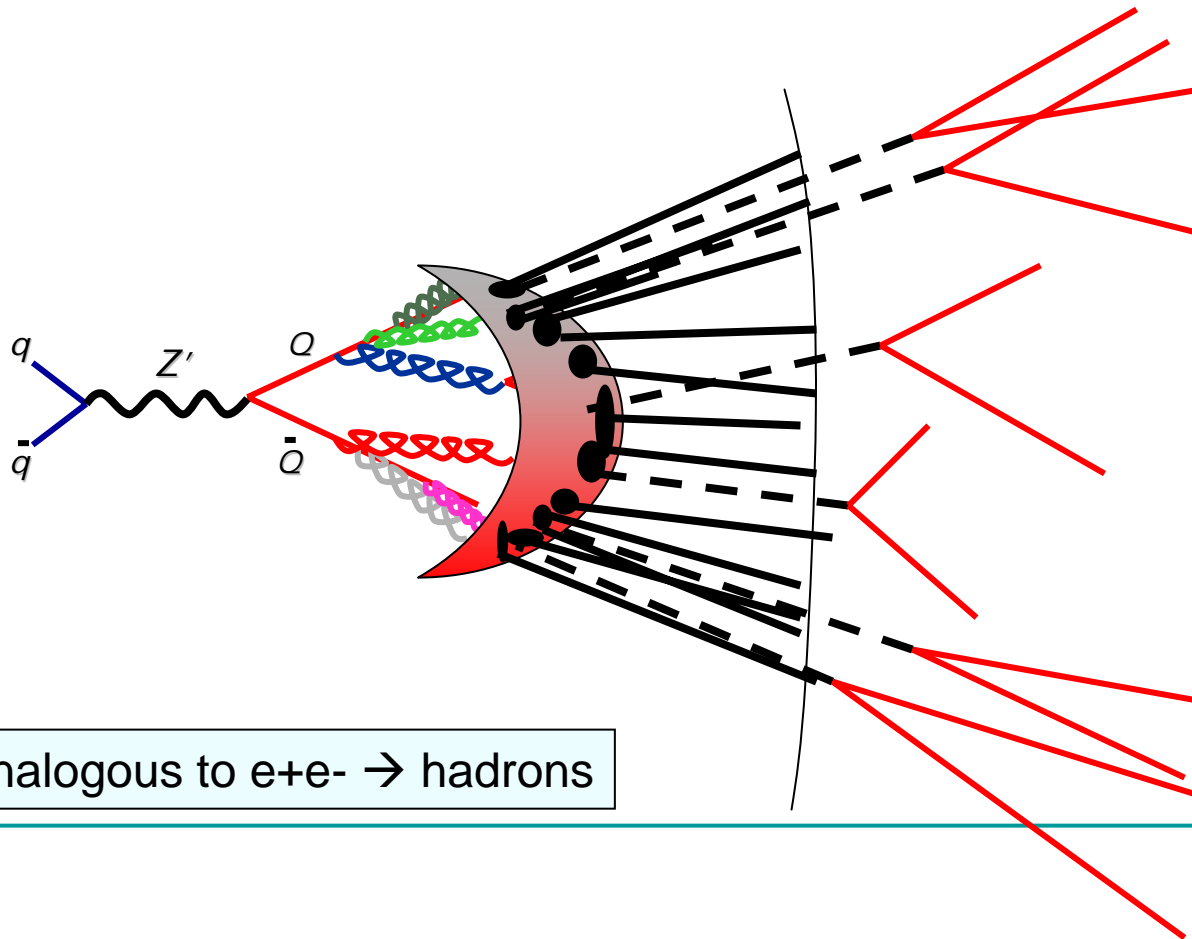


Analogous to $e^+e^- \rightarrow \text{hadrons}$

A Confining Hidden Valley

MJS + Zurek hep-ph/0604261

Some v -hadrons are stable and therefore invisible

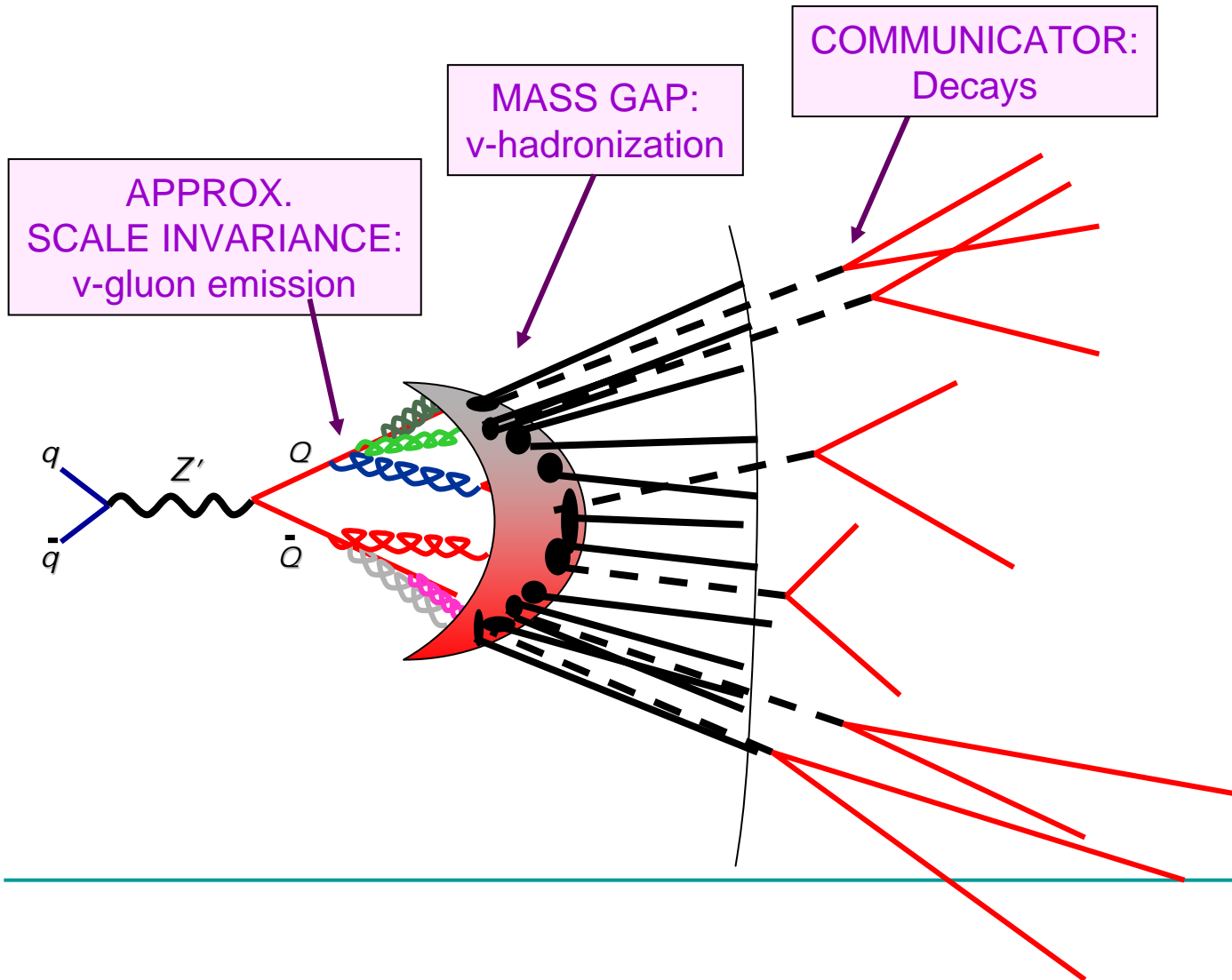


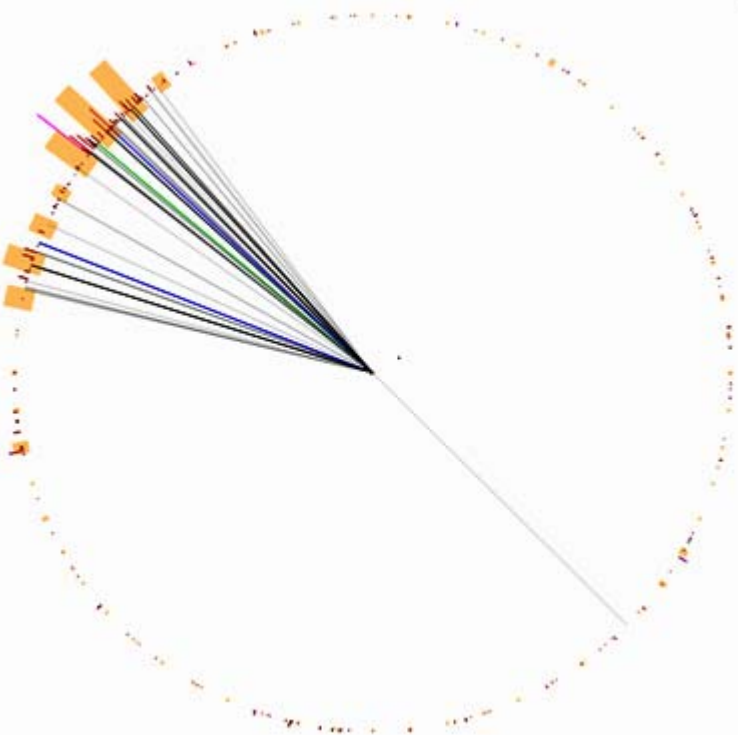
But some v -hadrons decay in the detector to visible particles, such as $b\bar{b}$ pairs, $q\bar{q}$ pairs, leptons etc.

Analogous to $e^+e^- \rightarrow \text{hadrons}$

High Multiplicity?

MJS + Zurek hep-ph/0604261





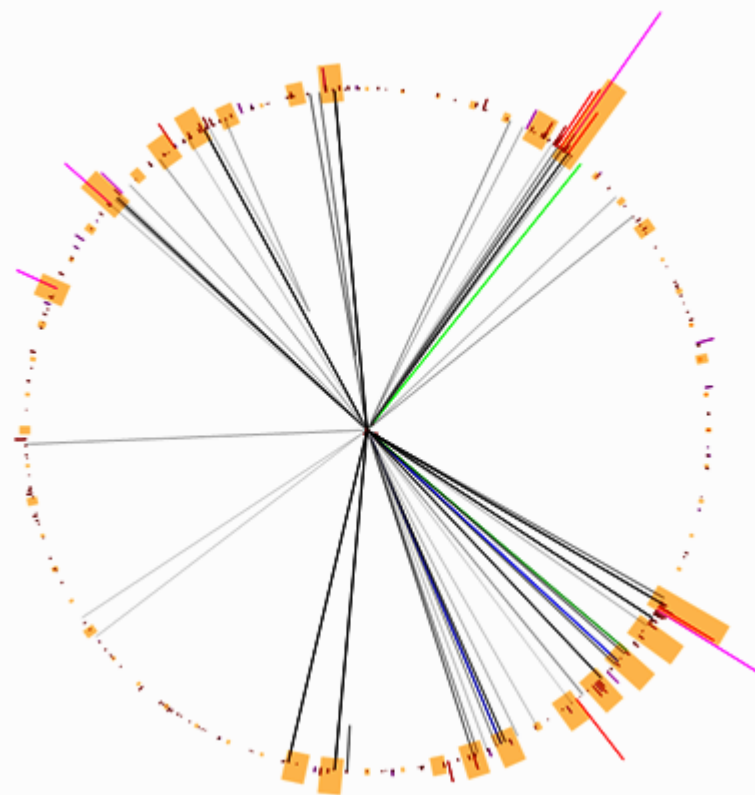
Average: 8 b's
Max: 22 b's

MJS, in preparation

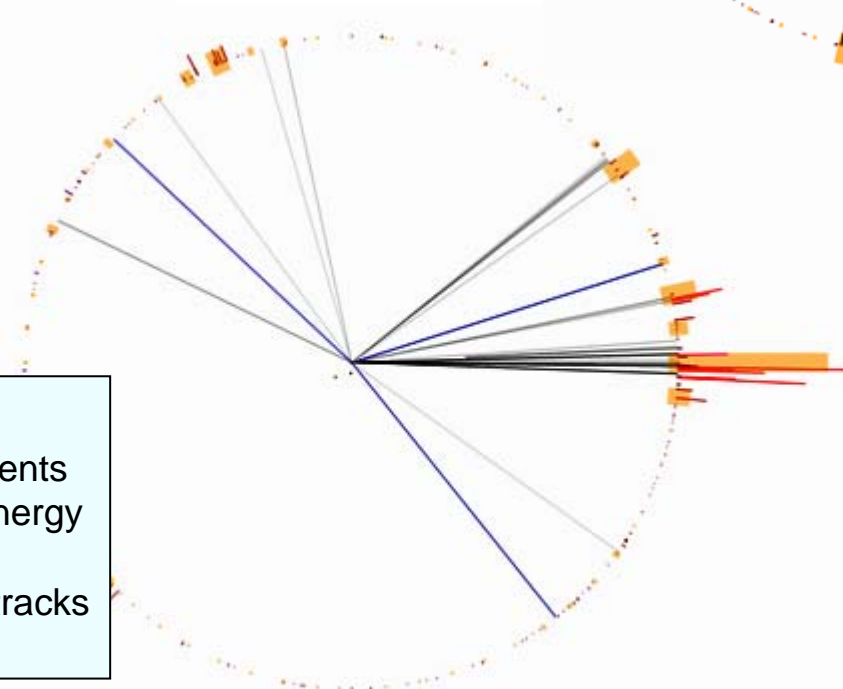
Search Strategy:

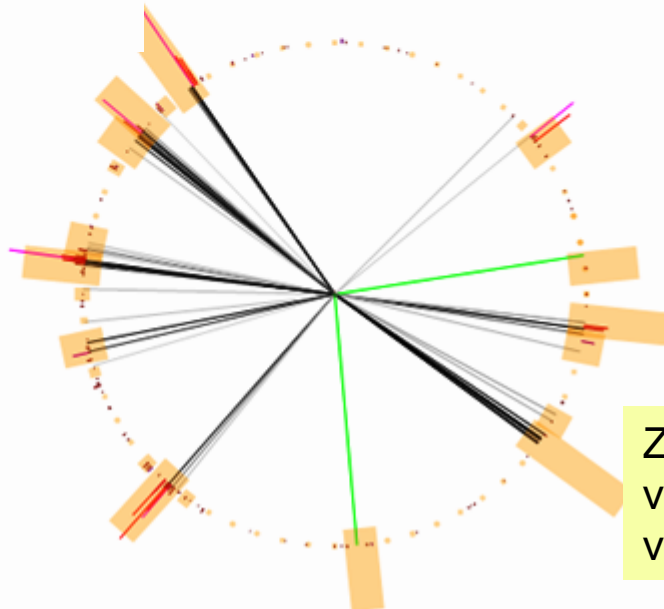
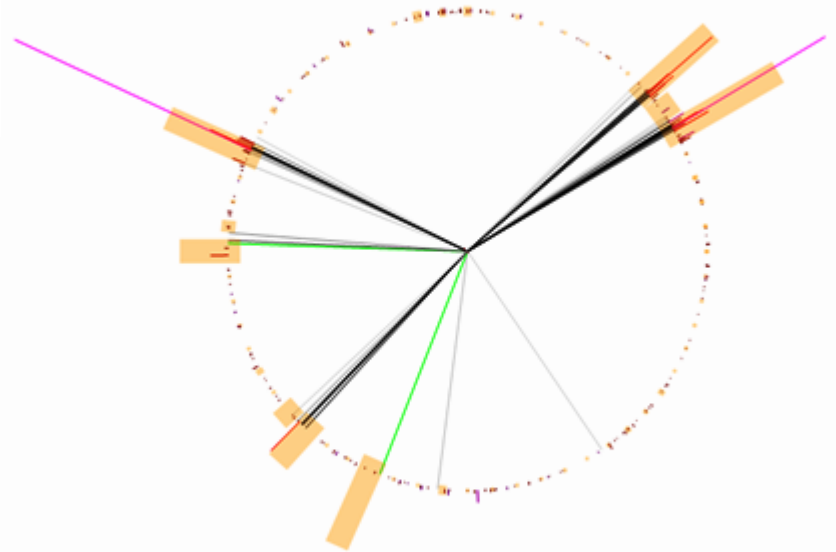
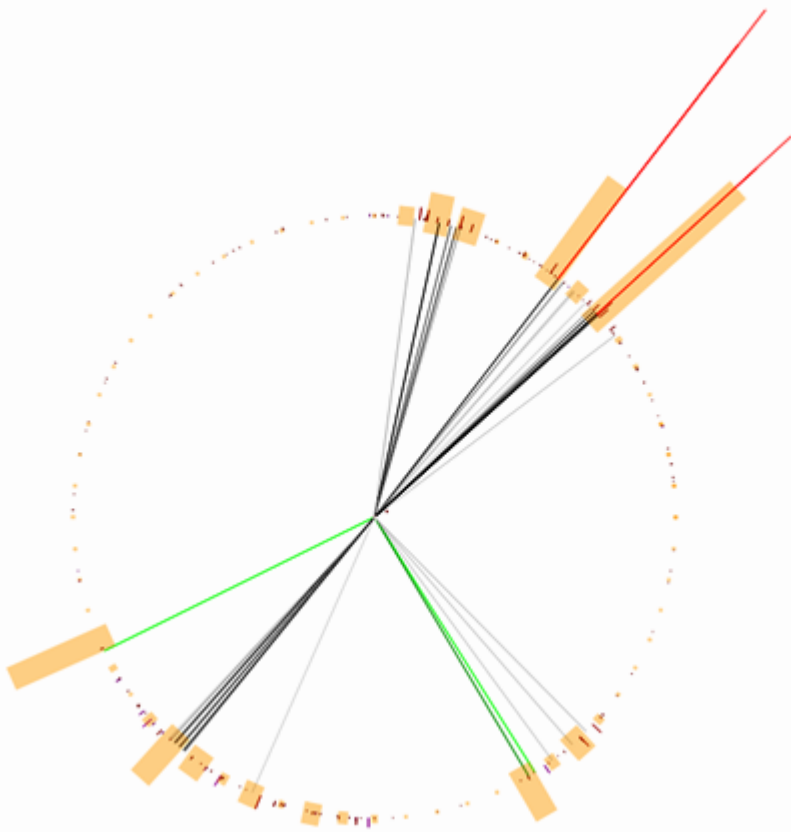
Find di-jet resonance in busy events

- missing and total transverse energy
- # jets
- # B meson vertices/displaced tracks
- Single jet invariant mass ?



Z' mass = 3.2 TeV
 ν -pi mass = 50 GeV
Flavor-off-diagonal
 ν -pions **stable**





Search Strategy:

Find dilepton resonance in busy events

- Event Shapes
- Dilepton invariant mass

Han, Si, Zurek + MJS,
2007

Z' mass = 3 TeV
v-eta' mass = 150 GeV
v-omega mass = 210 GeV

- If long-lived, light ν -hadrons
 - Spectacular events
 - But constraints from LEP \rightarrow rare
- **Must be detected with very high efficiency**
 - Online trigger to avoid discarding
 - Offline reconstruction to identify or at least flag



“Glass Elephant”

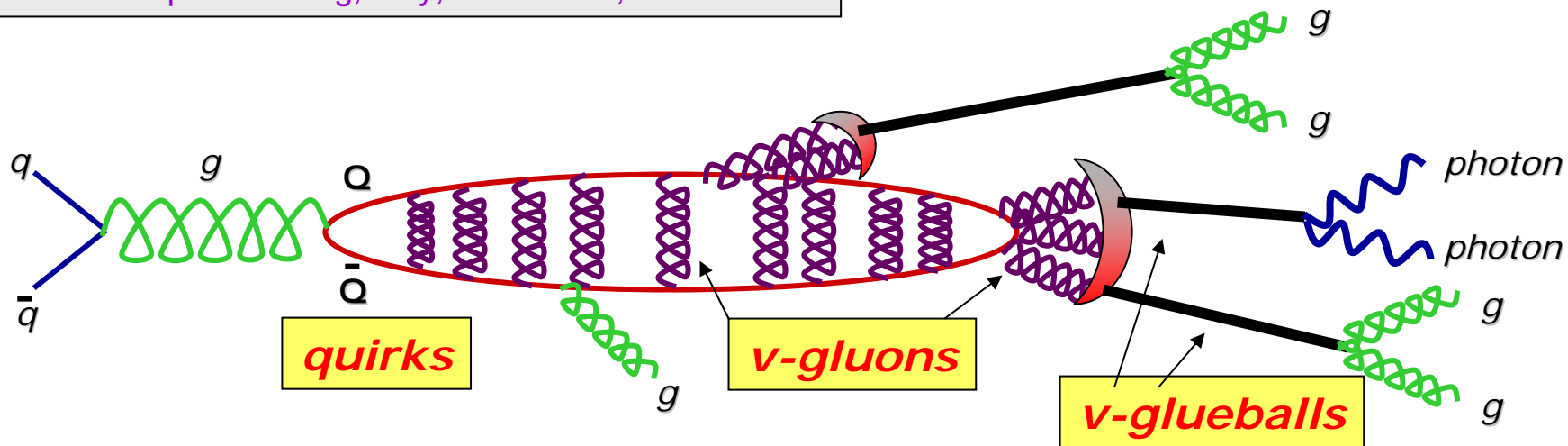
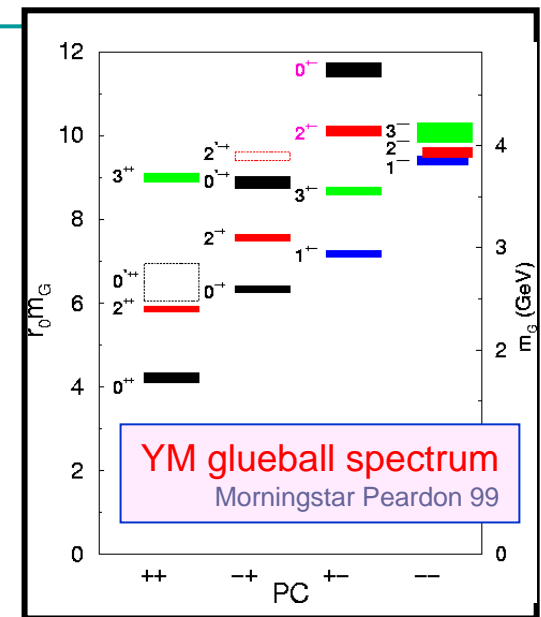
v-Glueballs

TeV-mass quirk production/annihilation

MJS + Zurek hep-ph/0604261
Juknevich, Melnikov, MJS in prep

- Quirk: Matter charged under SM and confining group...
- ...such that confining string cannot break

Lower-scale “quirks”: Kang, Luty, Nasri 2006; see next talk



Common Predictions of HV Scenario

- Higgs “Portal” allows higgs to mix into hidden sector

Recent work

Schabinger Wells [hep-ph/0509209](#)

MJS + Zurek [hep-ph/0604261](#)

Patt Wilczek [hep-ph/0605188](#)

- In Hidden Valley → new Higgs decays

MJS + Zurek [hep-ph/0604261](#), [hep-ph/0605193](#)

- $H \rightarrow XX$ with X decays **displaced** → new discovery mode

See also

Carpenter Kaplan Rhee 06

- $H \rightarrow XX$ with $X \rightarrow$ dileptons or dijets or diphotons

- $H \rightarrow XXX, XXXX$, etc

See also

Chang Fox Weiner 05

See also

Dermisek Gunion 04

Schabinger Wells 05

Chang Fox Weiner 06

Bowen Cui Wells 07

Gopalakrishna et al. 08

- Also true for other scalars

Common Predictions of HV Scenario

- Higgs “Portal” allows higgs to mix into hidden sector

Recent work
Schabinger Wells [hep-ph/0509209](#)
MJS + Zurek [hep-ph/0604261](#)
Patt Wilczek [hep-ph/0605188](#)

- In Hidden Valley → new Higgs decays

MJS + Zurek [hep-ph/0604261](#), [hep-ph/0605193](#)

- $H \rightarrow XX$ with X decays **displaced** → new discovery mode
- $H \rightarrow XX$ with $X \rightarrow$ dileptons or dijets or
- $H \rightarrow XXX, XXXX$, etc

See also
[Carpenter Kaplan Rhee 06](#)

See also
[Chang Fox Weiner 05](#)

- Also true for other scalars

CDF/D0: new searches carried out

Very difficult to trigger at ATLAS/CMS...
New ATLAS trigger strategy (internal note)

LHCb opportunity!!

Common Predictions of HV Scenario

- Big effect on SUSY

- LSP of SM sector can decay to the valley LSP...

MJS hep-ph/0607160

- ... plus SM particles ...
 - ... and/or plus $\tilde{\nu}$ -particles, which then decay to SM particles

- Either $\tilde{\nu}$ -particles or LSP may be long-lived

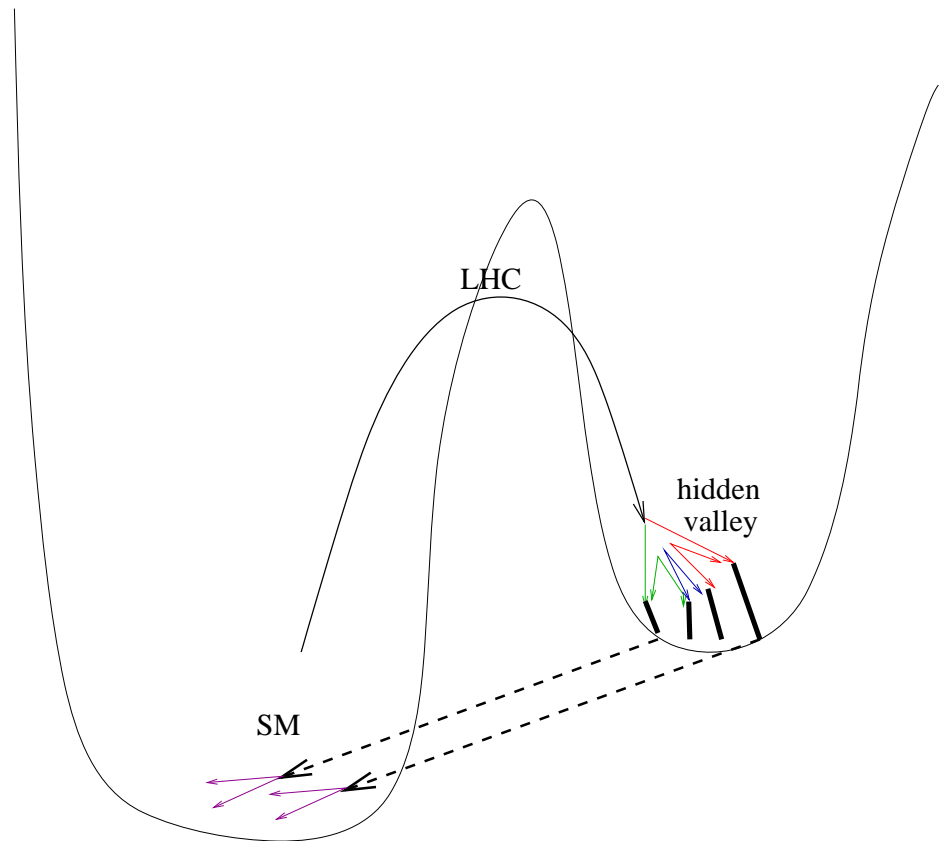
- Same effect applies to any theory w/ new global charge

- UED with KK parity
 - Little Higgs with non-anomalous T parity
 - Many other models

*Generalizes SUSY phenomenology from 90s
e.g. GMSB, Anomaly, etc.*

Remark on Unparticle Models

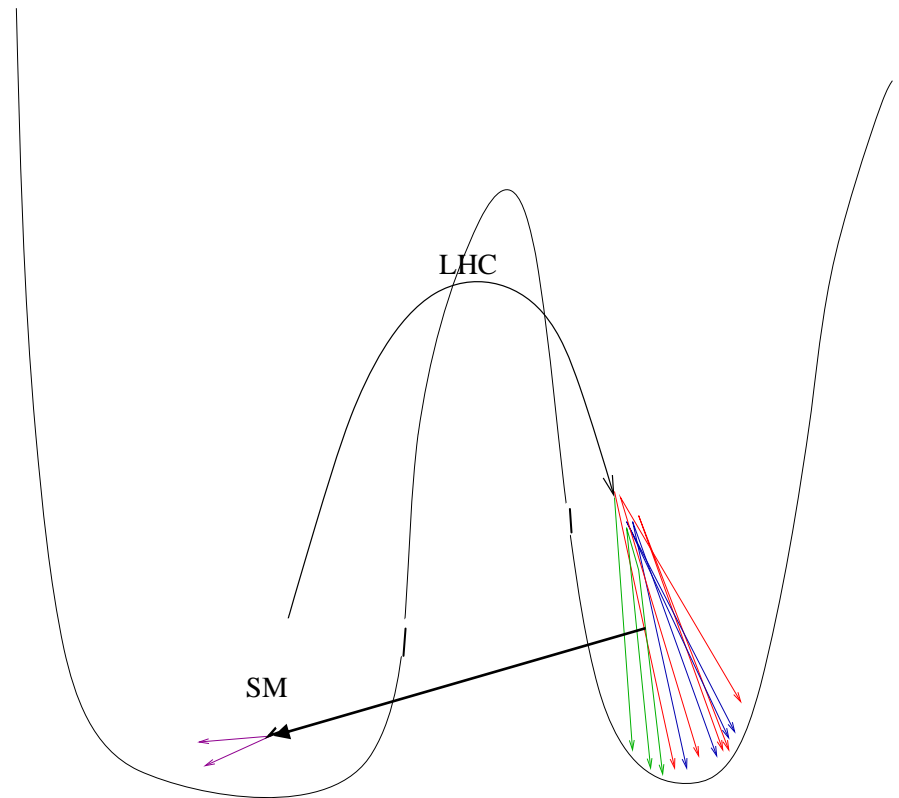
MJS-Zurek 2006



Remark on Unparticle Models

- In Unparticle models
 - a scale-invariant hidden sector generates indirect effects on observables
 - Events with MET
 - Rare virtual effects

Georgi 2007



Remark on Unparticle Models

- In Unparticle models

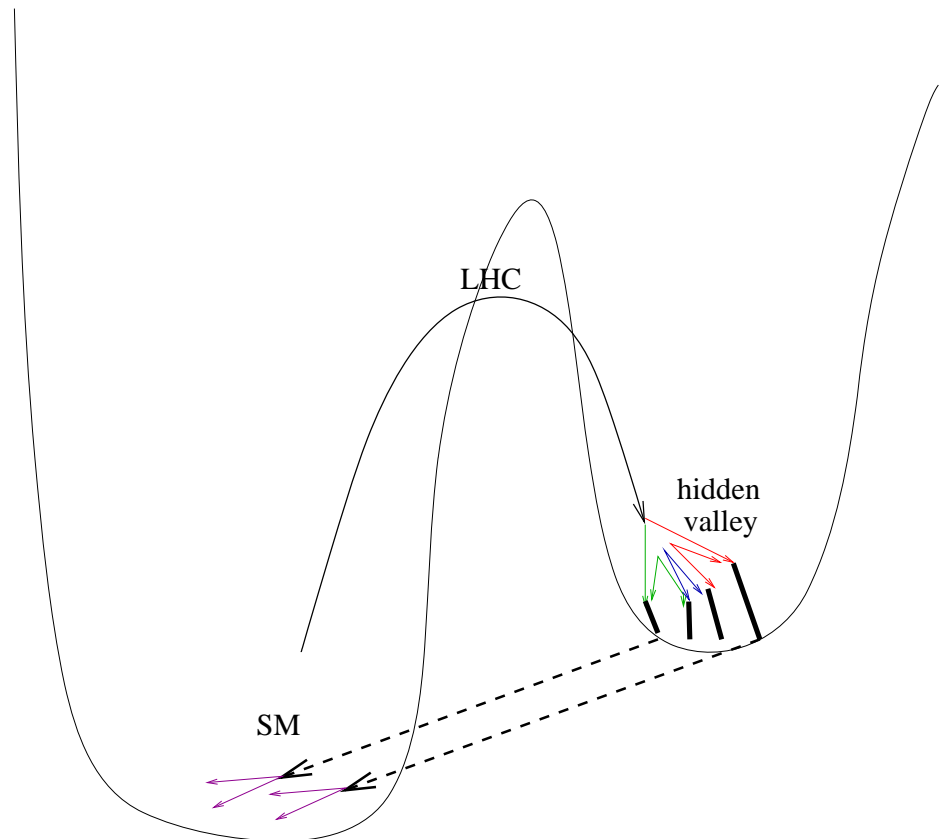
- a scale-invariant hidden sector generates indirect effects on observables
 - Events with MET
 - Rare virtual effects

- With large mass gap, model becomes a hidden valley

- Scale-symmetry breaking can lead to **direct, common, model-dependent, observable effects**

MJS 0801.

MJS-Zurek 2006

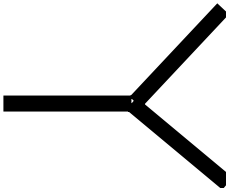


Unparticles vs. Conventional Language

- Unparticle \mathbf{U}

- Dimension D of \mathbf{U}


- Propagator 

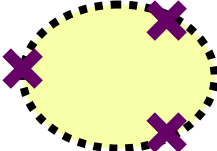
- Interactions 

- Couplings to SM 

- Gauge-Invariant Local Composite Operator $O(x)$

- Scaling Dimension D of $O(x)$

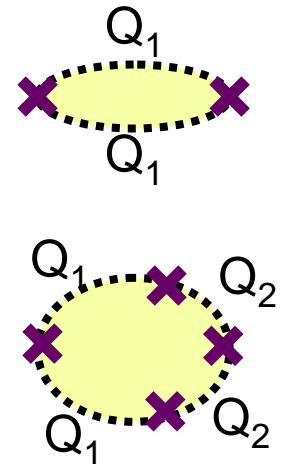
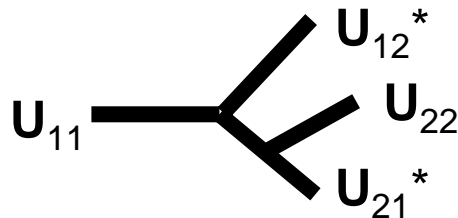
- 2-pt. Func $\langle O(x) O(y) \rangle$ 

- Higher-pt. Func.
 $\langle O(x) O(y) O(z) \rangle$ 

- Couplings to SM 

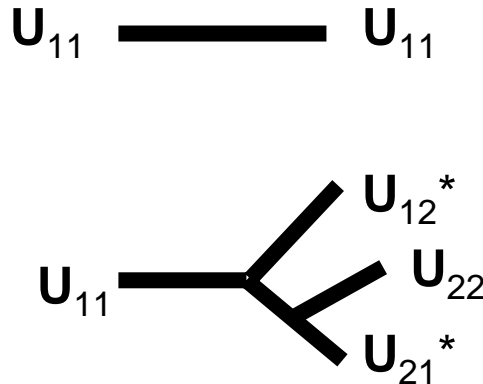
To Illustrate: A simple unparticle model

- Supersymmetric $SO(N)$ QCD with F flavors Q_j in hidden sector
- Conformal Fixed Point for $3(N-2) > F > \frac{3}{2}(N-2)$ Seiberg 1994
- Unparticles $\mathbf{U}_{JK} = Q_J Q_K$ have $D = 2 - \frac{3(F-N+2)}{2F}$ [thus $2 > D > 1$]

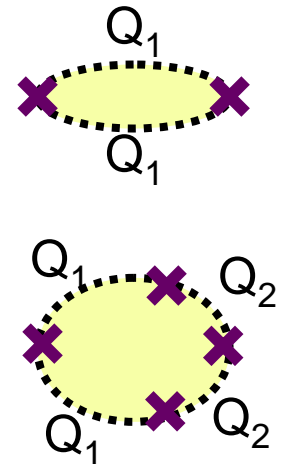


To Illustrate: A simple unparticle model

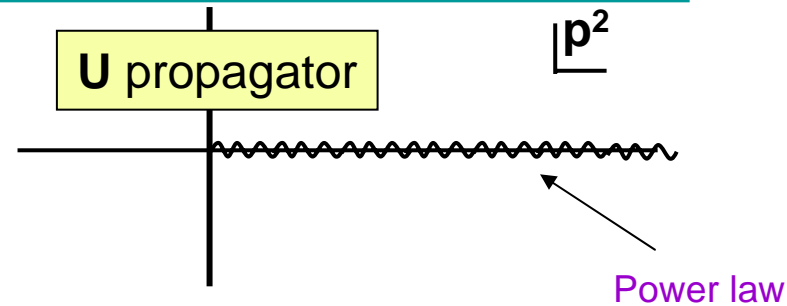
- Supersymmetric $SO(N)$ QCD with F flavors in hidden sector
- Conformal Fixed Point for $3(N-2) > F > \frac{3}{2}(N-2)$
- Unparticles $\mathbf{U}_{JK} = Q_J Q_K$ have $D = 2 - \frac{3(F-N+2)}{2F}$ [thus $2 > D > 1$]



As $F \rightarrow 3(N-2)$
 Fixed pt. cplg $\rightarrow 0$
 (Banks-Zaks limit)
 Q becomes free
 but
 U does not!



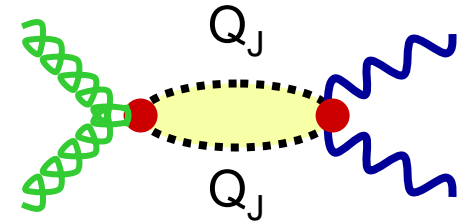
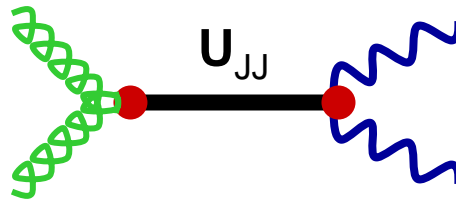
Scale-Invariant Sector



Couple to SM gauge bosons: $1/M^p (\sum c_{JK} \mathbf{U}_{JK}) \text{Tr} [F^2]$

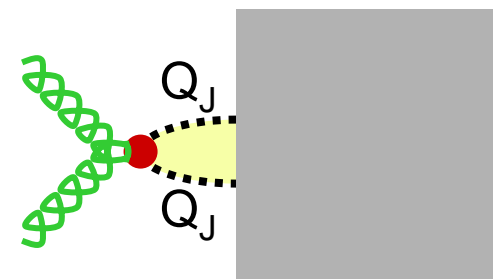
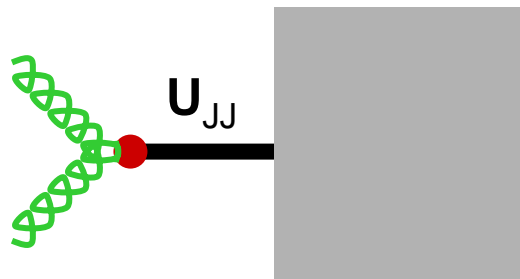
$g g \rightarrow 2 \text{ photons}$

$$\sim c^2 M^{-2p}$$



$g g \rightarrow \text{nothing}$

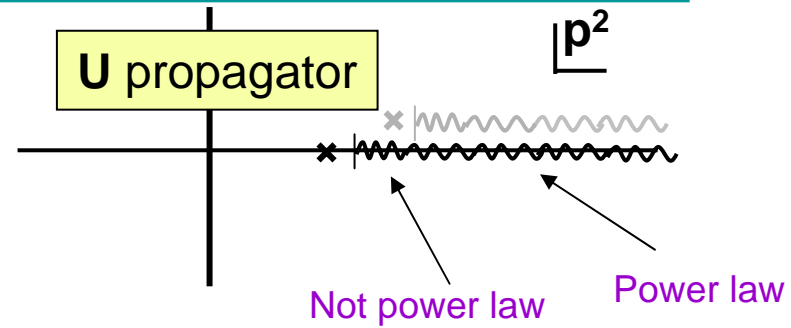
$$\sim c^1 M^{-p}$$



Dominant Process

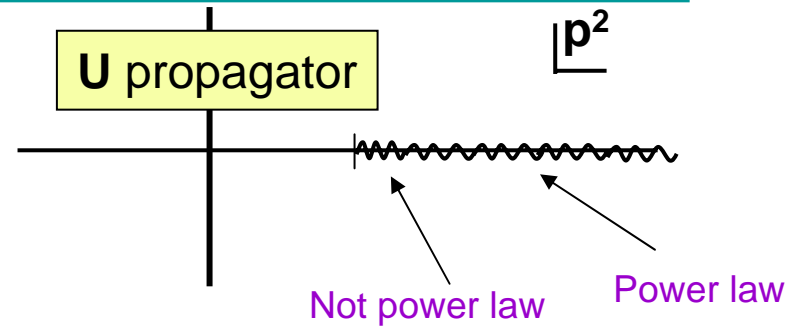
Masses m_{JK} for the Q_J

- Add term $m_{JK}^2 Q_J Q_K^* = m_{JK}^2 \mathbf{V}_{JK}$



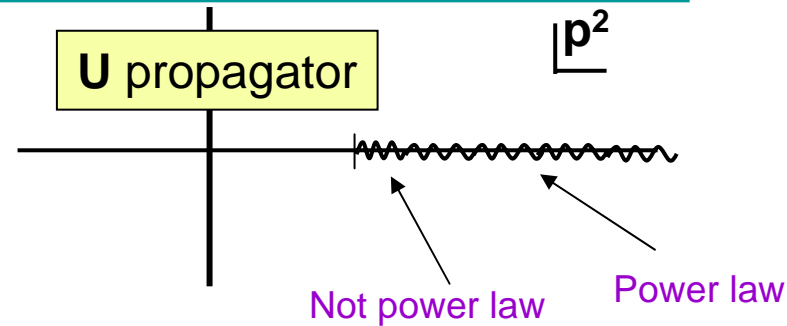
Masses m_{JK} for the Q_J

- Add term $m_{JK}^2 Q_J Q_K^* = m_{JK}^2 \mathbf{V}_{JK}$



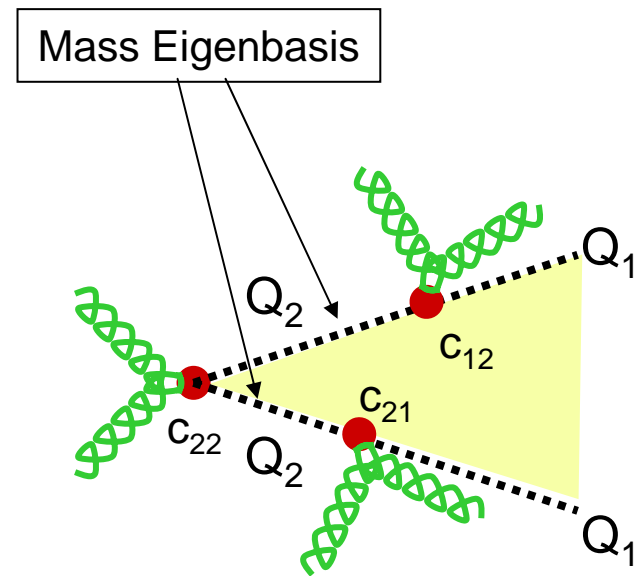
Masses m_{JK} for the Q_J

- Add term $m_{JK}^2 Q_J Q_K^* = m_{JK}^2 \mathbf{V}_{JK}$
- Suppose m_{JK}^2 , c_{JK} not simultaneously diagonal



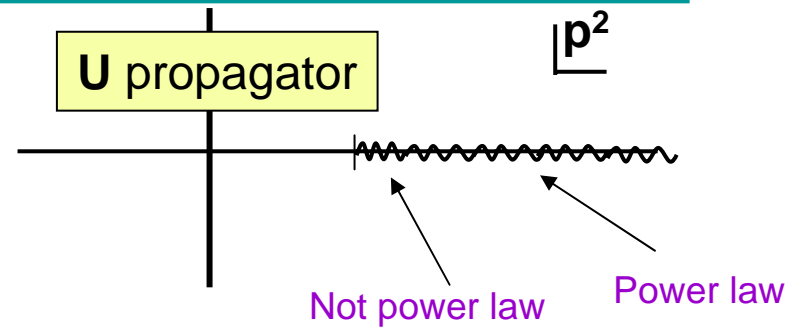
$g g \rightarrow 4 \text{ jets} + \text{MET}$

$\sim c_{22}^1 M^{-p}$ **[NOT $c^3 M^{-3p}$!]**



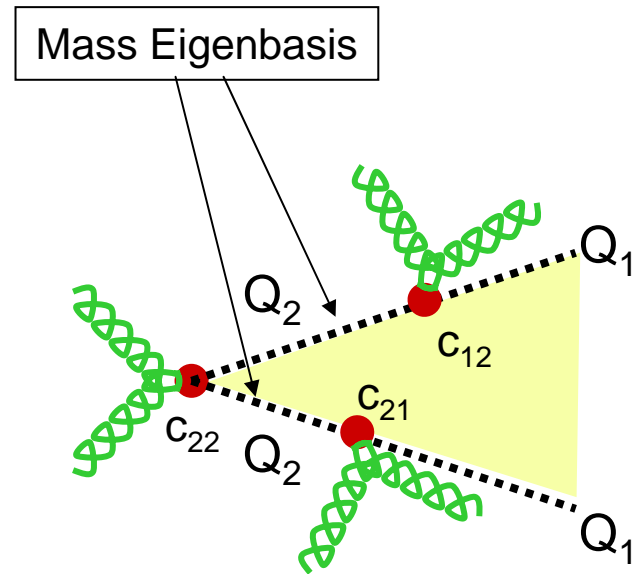
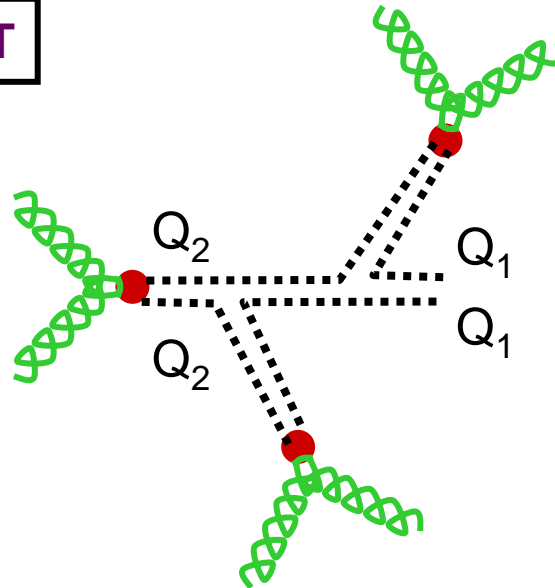
Masses m_{JK} for the Q_J

- Add term $m_{JK}^2 Q_J Q_K^* = m_{JK}^2 \mathbf{V}_{JK}$
- Suppose m_{JK}^2 , c_{JK} not simultaneously diagonal



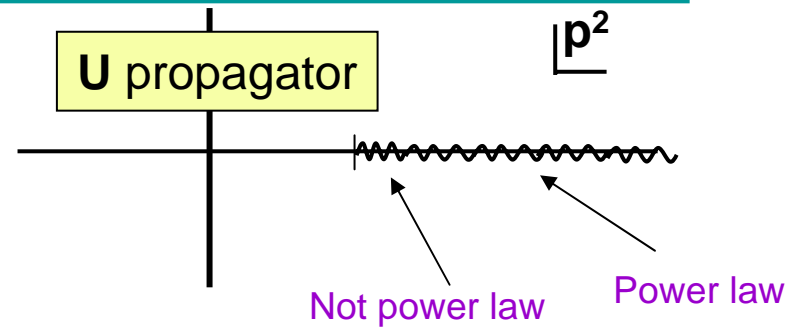
$g g \rightarrow 4 \text{ jets} + \text{MET}$

$\sim c_{22}^1 M^{-p}$



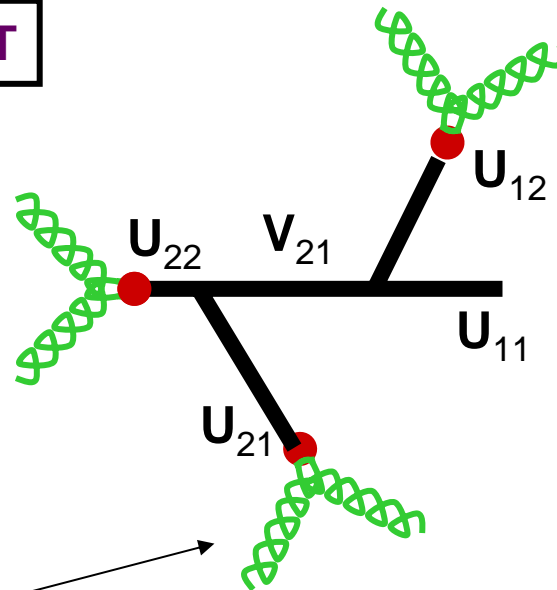
Masses m_{JK} for the Q_J

- Add term $m_{JK}^2 Q_J Q_K^* = m_{JK}^2 V_{JK}$
- Suppose m_{JK}^2 , c_{JK} not simultaneously diagonal



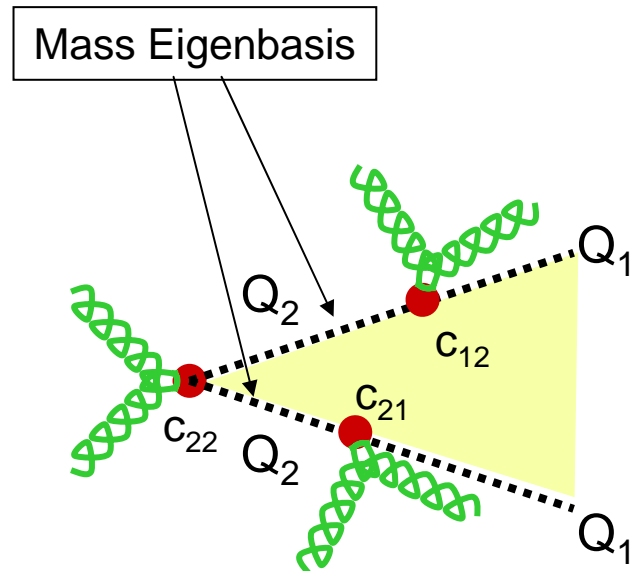
$g g \rightarrow 4 \text{ jets} + \text{MET}$

$\sim c_{22}^1 M^{-p}$



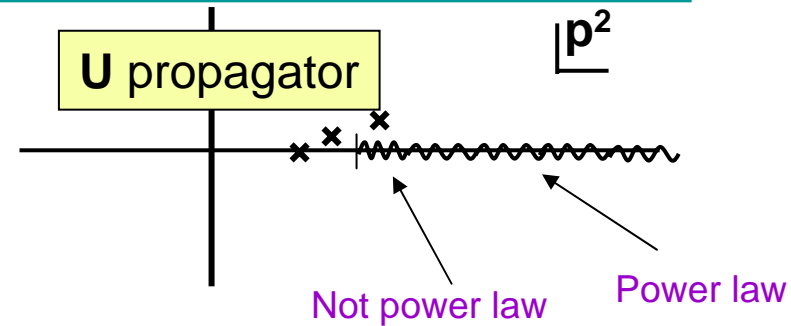
$$V_{21} = Q_2 Q_1^*$$

Below the cut



Hidden sector is not invisible
Visible production amplitude is order(c)
 Higher-multiplicity
 Can't use D to predict kinematics

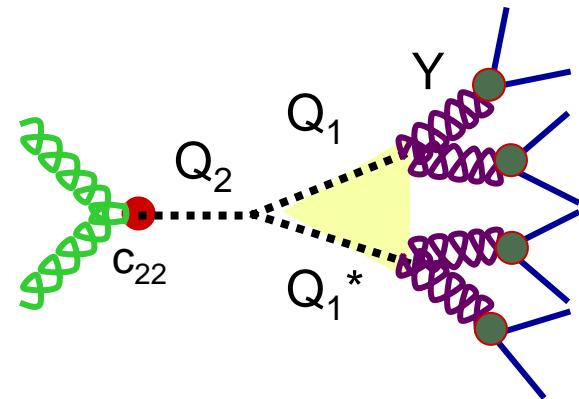
$$m_J \text{ and } \langle Q_J \rangle = v_J$$



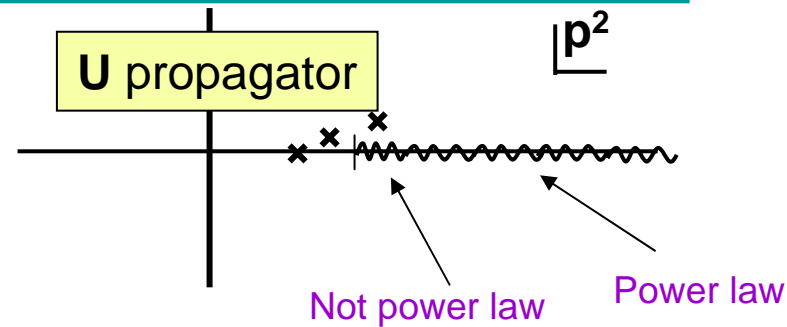
- SUSY potential now has cubic terms: $Q_J \rightarrow Q_K Q_K^*$
- Hidden massive v-gluons Y
- If Y mixes with Z boson, then $Y \rightarrow \text{SM fermions}$

$g g \rightarrow 8 \text{ quarks/leptons}$

$$\sim c_{22}^1 M^{-p}$$



$$m_J \text{ and } \langle Q_J \rangle = v_J$$

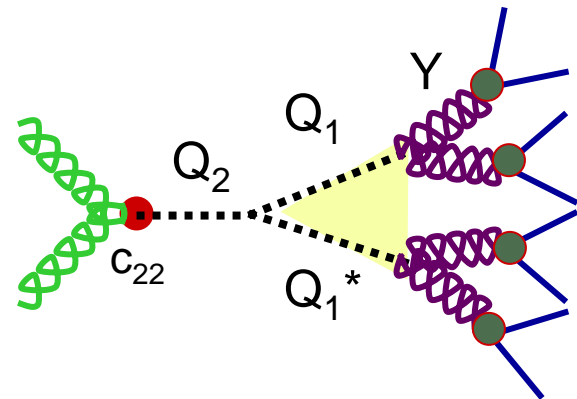
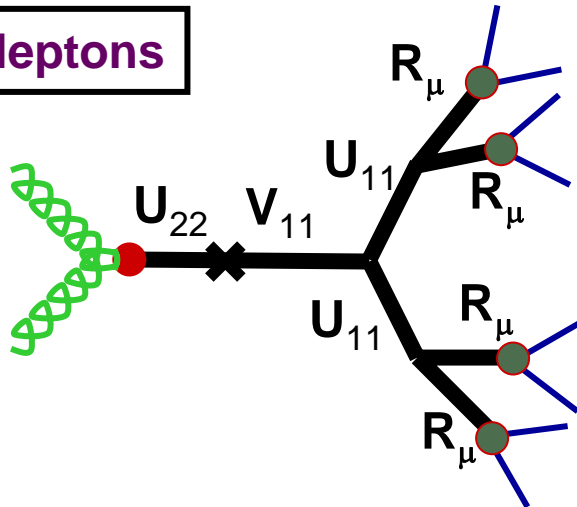


- Shifted U: linear + quadratic in shifted Q
- U's interact with unparticle $R_\mu = Q^* D_\mu Q$
- If Z boson couples to $R_\mu \dots$

- SUSY potential now has cubic terms: $Q_J \rightarrow Q_K Q_K^*$
- Hidden massive v-gluons Y
- If Y mixes with Z boson, then $Y \rightarrow \text{SM fermions}$

$g g \rightarrow 8 \text{ quarks/leptons}$

$$\sim c_{22}^1 M^{-p}$$



Hidden sector is not invisible
Visible production amplitude is order(c)
 High-multiplicity, light neutral resonances
Possible displaced vertices

Lesson

- We cannot predict production rates, final states, or observables unless we know **precisely**
 - The nature of the hidden sector and its multipoint functions
 - The couplings of SM to that sector
 - The details of the breaking of scale invariance
- With mass gap, often D of unparticle cannot be measured at LHC

Example: $t \rightarrow c \mathbf{U}$: measure t/c relative kinematics to determine D

Signatures

- Gap $\ll \ll \ll m_{\text{top}}$: \mathbf{U} invisible \rightarrow Missing energy
- Gap $\ll m_{\text{top}}$: $\mathbf{U} \rightarrow$ many SM particles; \mathbf{U} daughters long-lived
- Gap $\ll m_{\text{top}}$: $\mathbf{U} \rightarrow$ many SM particles; \mathbf{U} daughters short-lived
- Gap not $\ll m_{\text{top}}$: $\mathbf{U} \rightarrow$ SM particles; t/c kinematics not determined by D

<u>HV</u>	<u>Un</u>
No	Yes
Yes	Yes
Yes	<u>Can't Observe</u>
Yes	No

U_n/HV and AdS/CFT

U propagator

$|p^2$

Stephanov Narrow-Tower Model [not “deconstruction”]

- AdS/CFT: mass gap turns $\langle U U \rangle$ into sum of poles
- Predicts tower of long-lived states (lifetimes $\rightarrow D$)



Un/HV and AdS/CFT

Stephanov Narrow-Tower Model [not “deconstruction”]

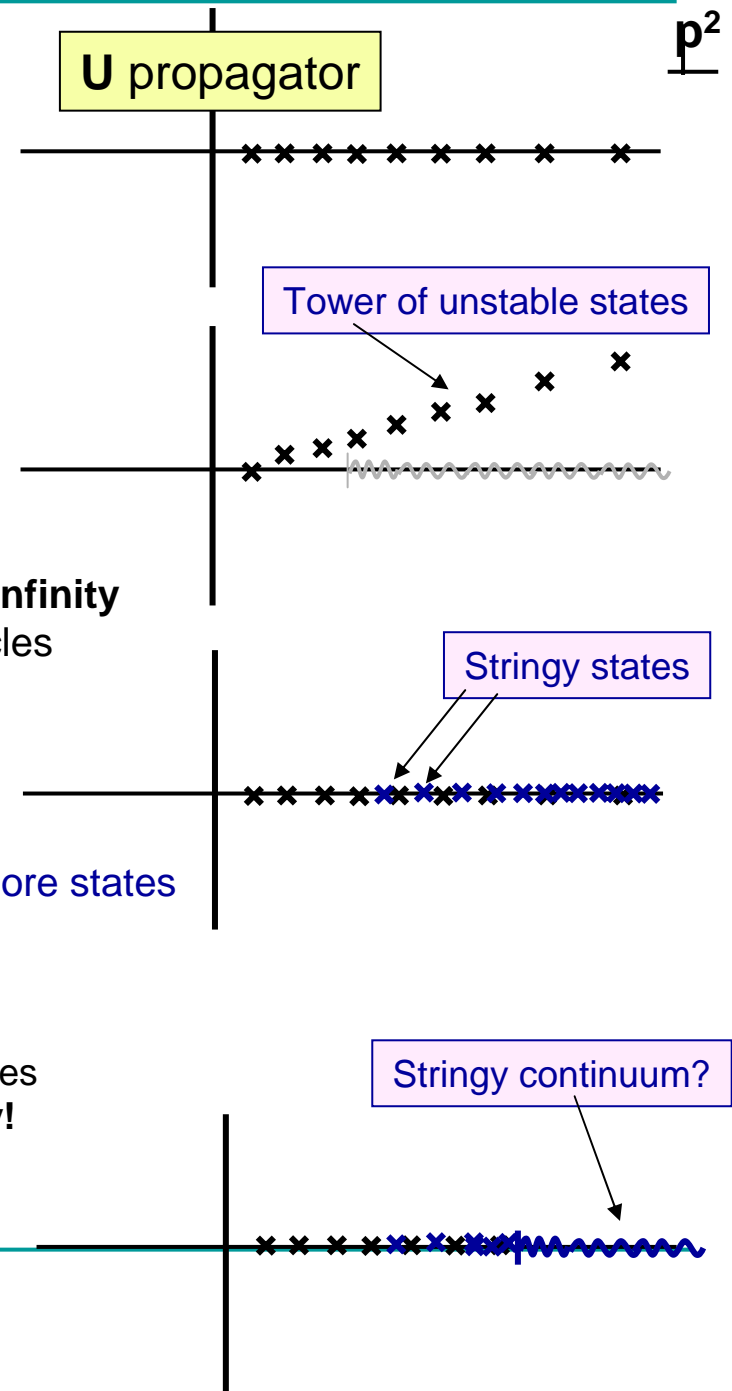
- AdS/CFT: mass gap turns $\langle \mathbf{U} \mathbf{U} \rangle$ into sum of poles
- Predicts tower of long-lived states (lifetimes $\rightarrow D$)

But neglects unparticle interactions...

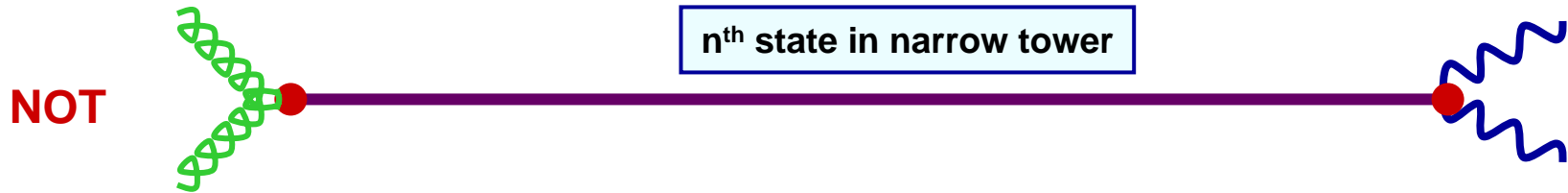
- AdS/CFT: tower = narrow hadronic resonances at $N \rightarrow \text{infinity}$
- **Finite** N : unparticle interacts with self and other unparticles
- Narrow states are wide! Also new cuts!

Also neglects strings...

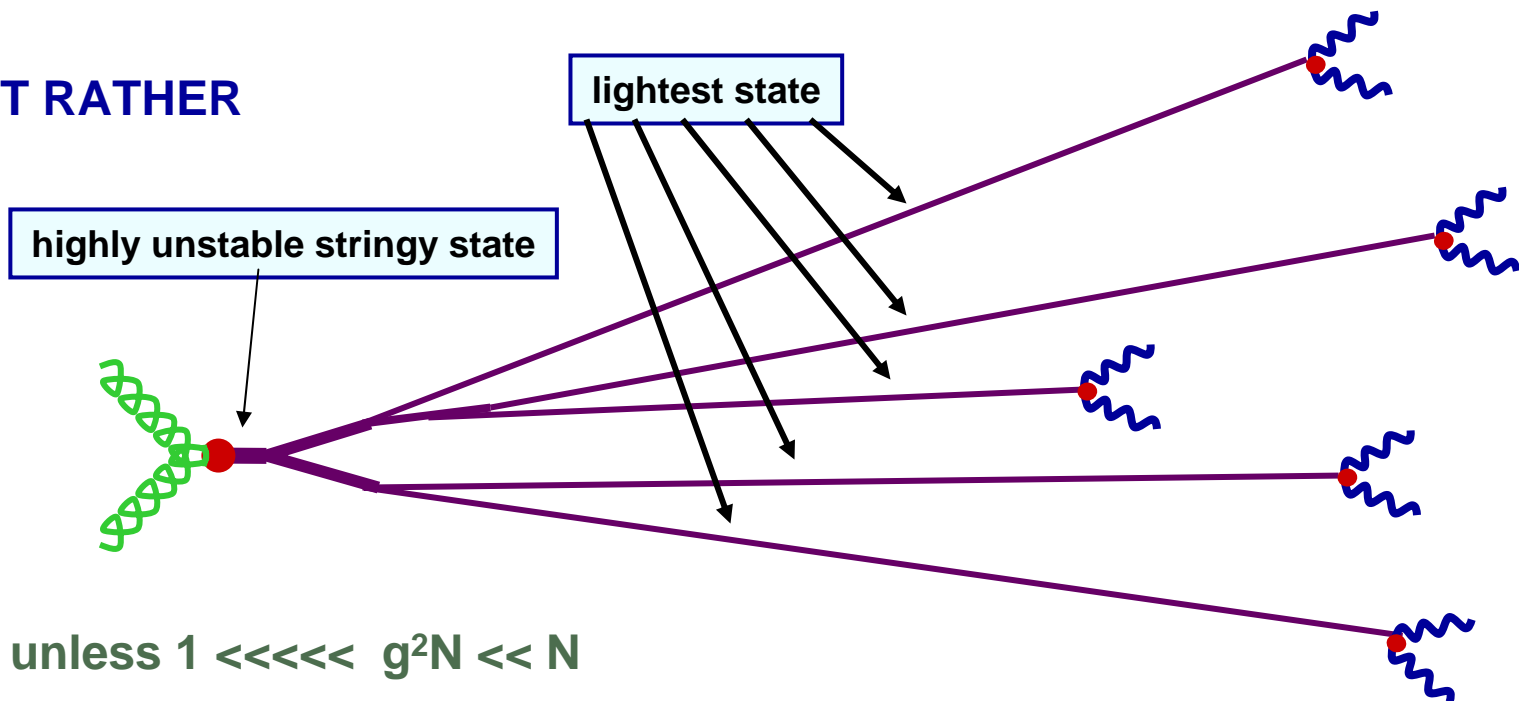
- At finite 't Hooft coupling g^2N , even at infinite N , **many more states**
- There may even be a **continuum** at infinite N , finite g^2N
 - D7 brane: quarkonium states plus unconfined quark states
 - **The continuum cannot be seen without string theory!**



Un/HV and AdS/CFT



BUT RATHER



Thus AdS/CFT DOES resemble QCD, QCD-like Hidden Valley

Caution on AdS/CFT usage

Sometimes, “ \sim ” is not enough

paraphrasing Matthias Neubert

For Analytic Structure of Correlation Functions,
“ $=$ ” is necessary

- In general, if one neglects
 - $1/N$ corrections
(bulk interactions, D-branes, metric back-reaction,...)
 - $1/(g^2N)$ corrections
(bulk classical string corrections, massive states, high-spin operators,...)

one might miss subtle and crucial details of the physical amplitudes

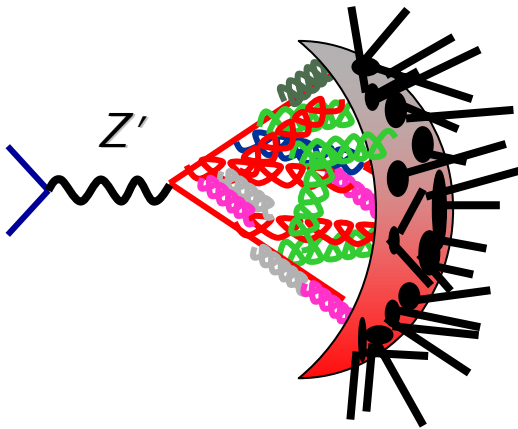
Predictions of hidden strong dynamics

Powerful gluon emission in scale-invariant region at strong 't Hooft coupling –

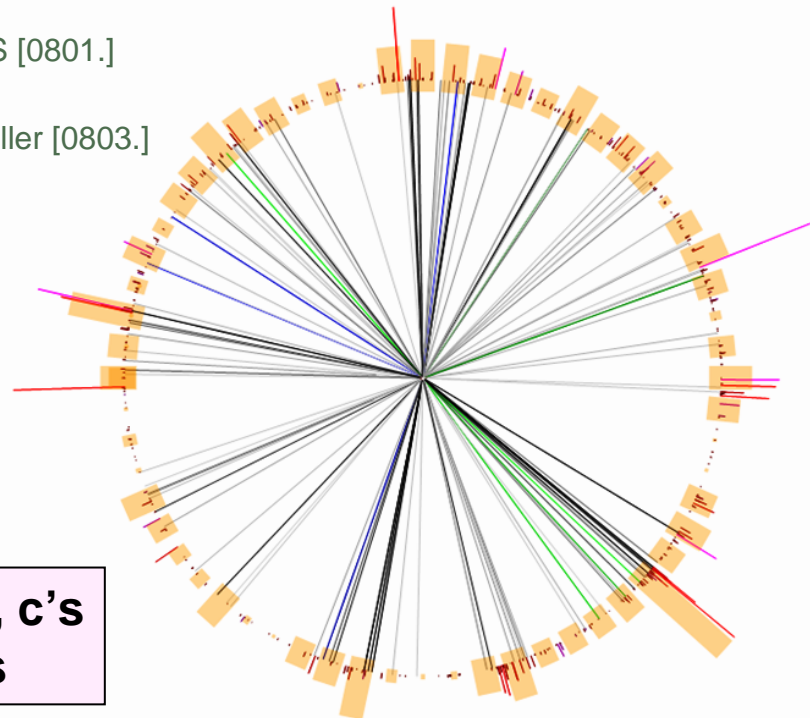
- from large anom. dims. of **high-spin high-dimension** operators

Predicts near-spherical high-multiplicity events

- Conjecture based on AdS/CFT arguments MJS [0801.]
- Recently proven – Hoffman/Maldacena [0803.]
see also conjecture in Hatta, Iancu, Mueller [0803.]



~ 54 b's, c's
~ 6 tau's



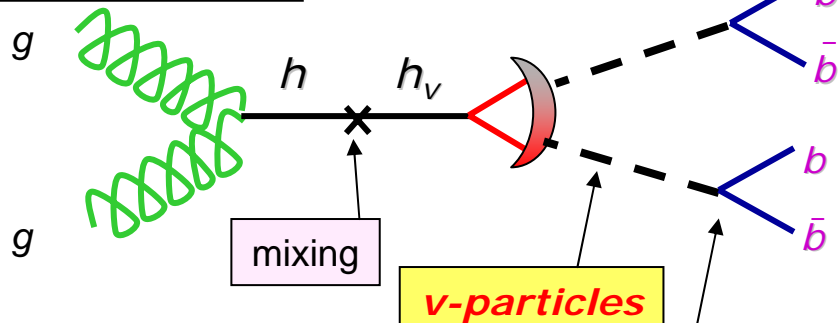
Summary

- Hidden valleys → rich and novel phenomena at LHC and beyond
 - High multiplicity events
 - Long-lived particles
 - Novel decays for Higgs, LSP/LTP/LKP, Z' , Z , etc.
- “Unparticle” models with mass gaps become HV models
 - Hidden sector becomes visible, at leading order in coupling
 - Unparticle interactions must be included
 - Must account for multiple Unparticles in any model
 - Unparticle dimension does not determine most observables
 - Many processes give classic Hidden Valley signatures
- AdS/CFT must be used with great care when making predictions
- For large g^2N hidden valley, AdS/CFT predicts:
 - Ultra-high-multiplicity quasi-spherical events with low- p_T jets/leptons

Backup Slides

MJS + Zurek
hep-ph/0604261
hep-ph/0605193

Displaced vertex



CDF/D0: new searches carried out

Very difficult to trigger at ATLAS/CMS...
New ATLAS trigger strategy (internal note)

LHCb opportunity!!

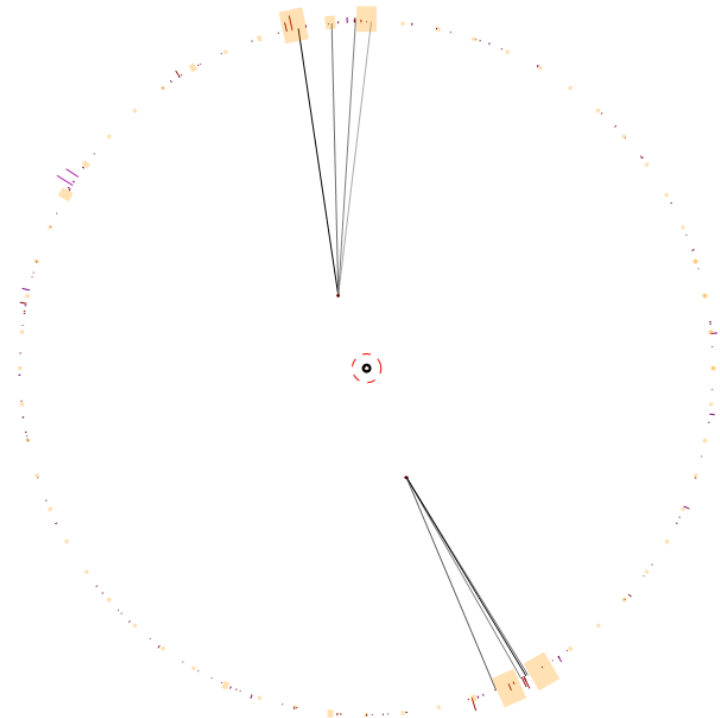
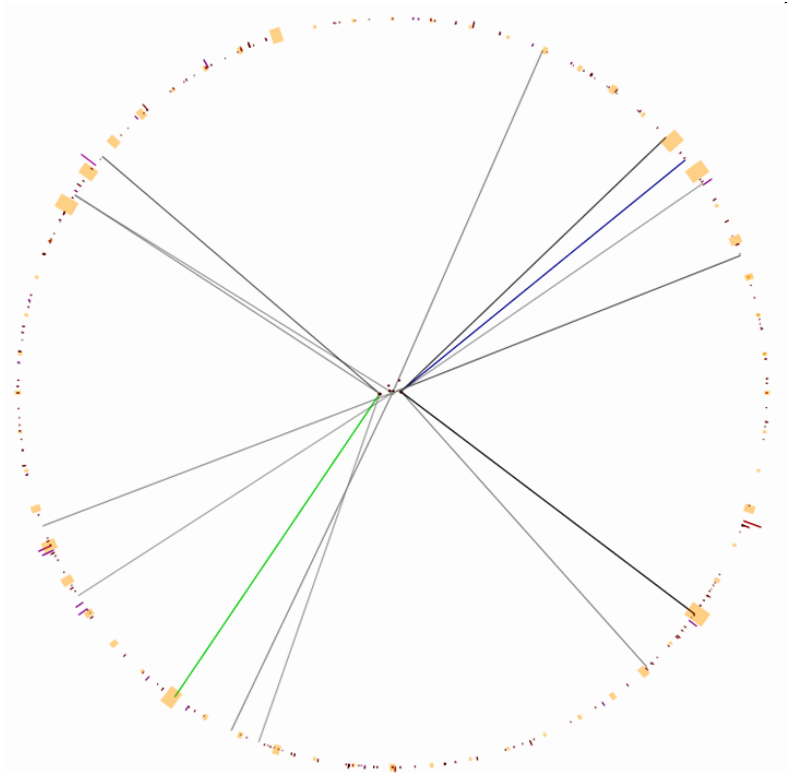
Similar Observations:

hep-ph/0607204 : Carpenter, Kaplan and Rhee

Precursor (LEP focus):

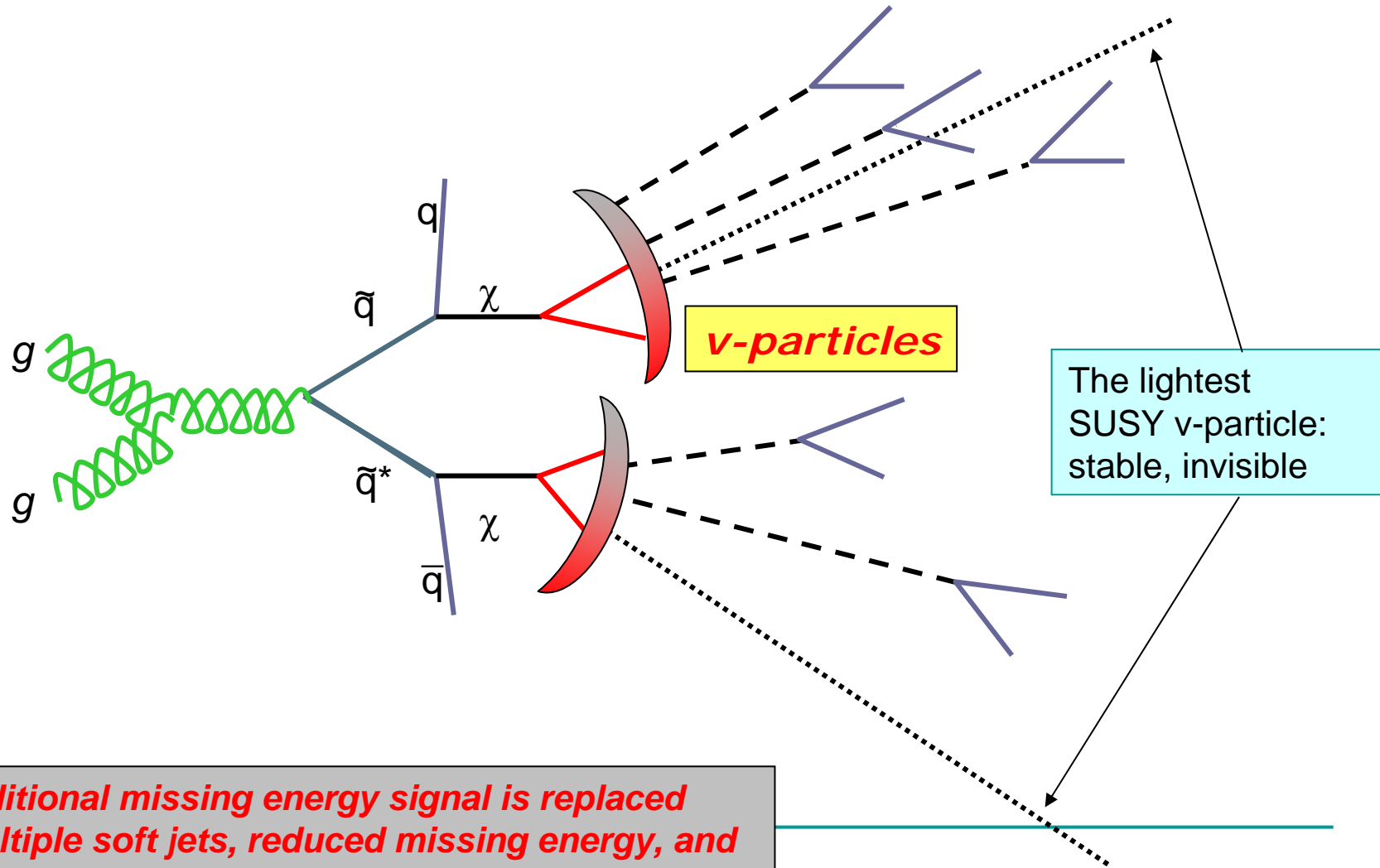
Chang, Fox, Weiner, limit of model in hep-ph/0511250

Displaced vertex

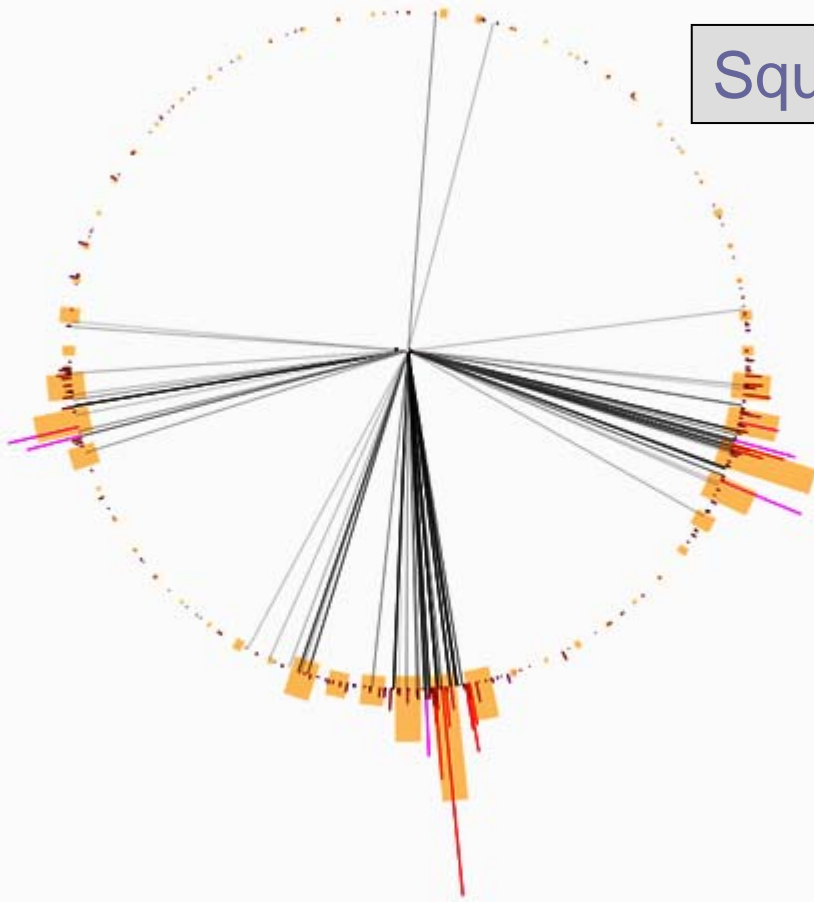


SUSY decays to the v -sector

MJS hep-ph/0607160

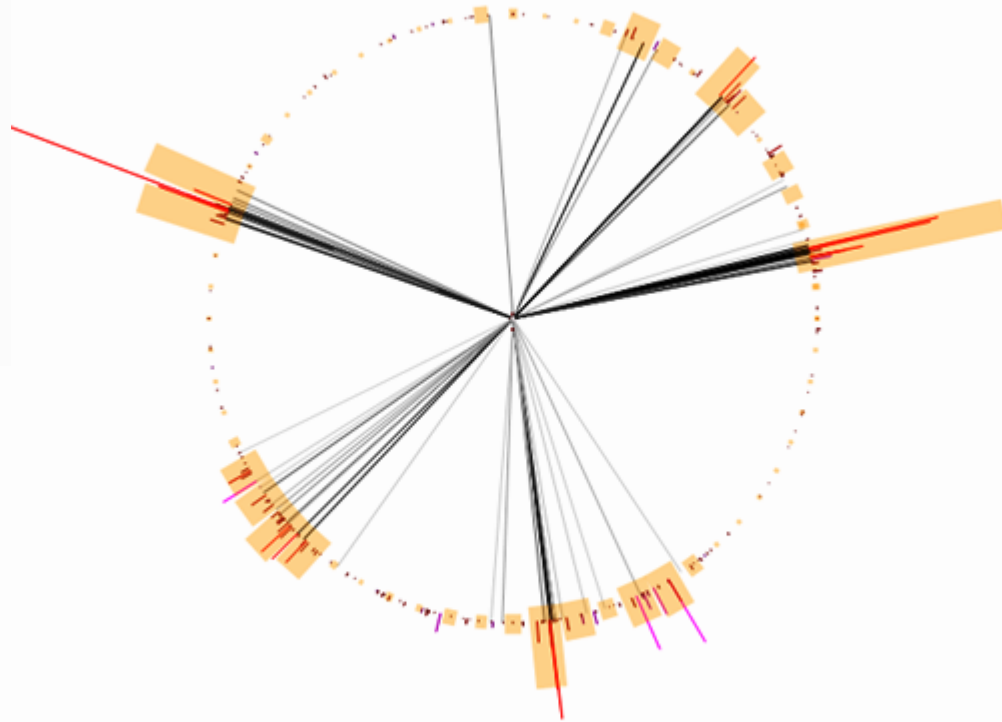


Squark-Antisquark Production at LHC



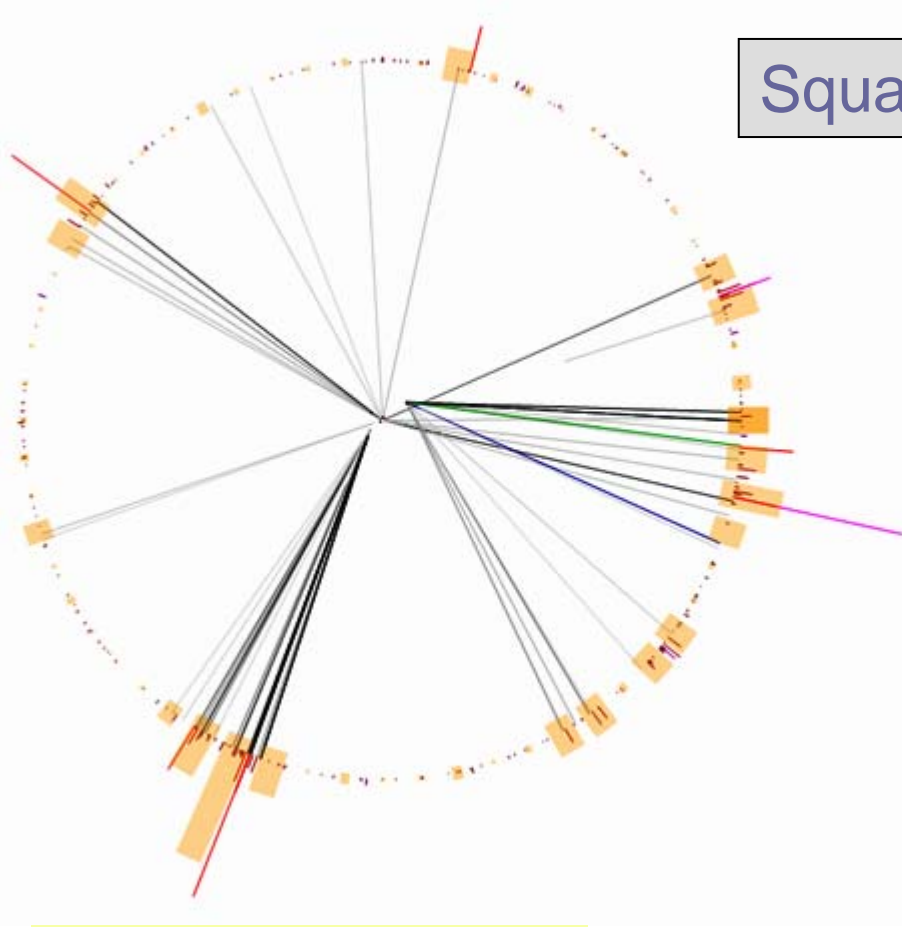
Stable Neutralino

Unstable Neutralino
Decaying to ν -Sector



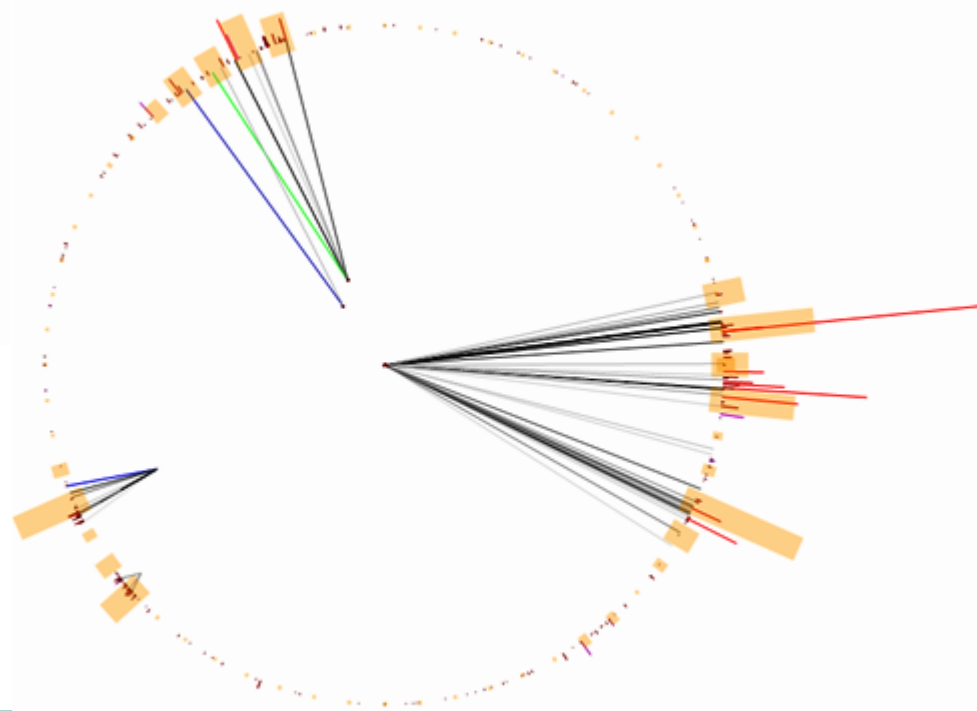
Hacked simulation using
Hidden Valley Monte Carlo 1.0
Mrenna, Skands and MJS

Squark-Antisquark Production at LHC



Long-Lived Neutralino
Prompt ν -Particle Decay

Prompt Neutralino Decay
Long-Lived ν -Particles

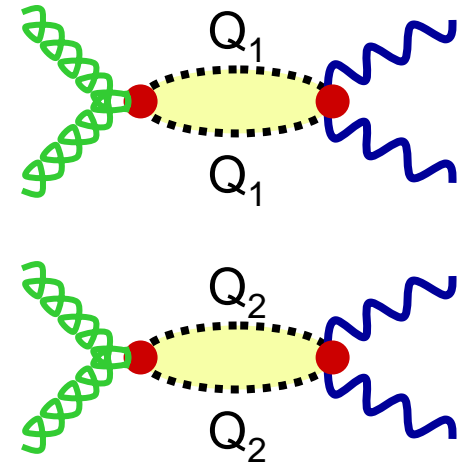
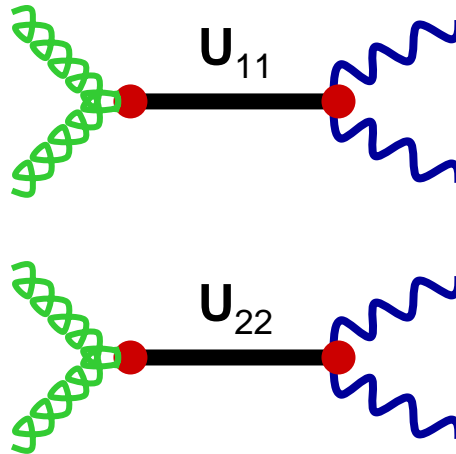


Hacked simulation using
Hidden Valley Monte Carlo 1.0
Mrenna, Skands and MJS

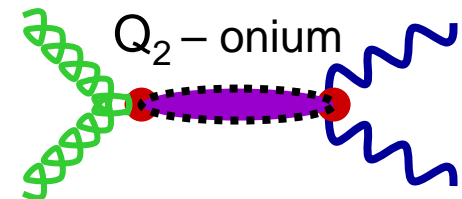
Masses m_J for the Q_J

$g g \rightarrow 2 \text{ photons}$

$$\sim c_J M^{-p}$$



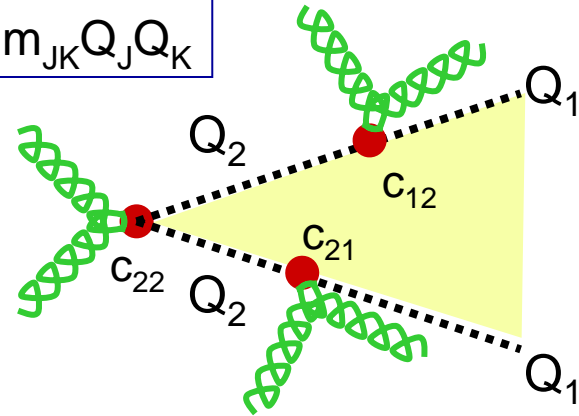
- Thresholds at energy = $2 m_1, 2 m_2, \dots$
- v -Quarkonium resonances below threshold
- Details depend on N, F, m_J



Hidden sector effects not predictable
Order- c^2 terms dominate observables

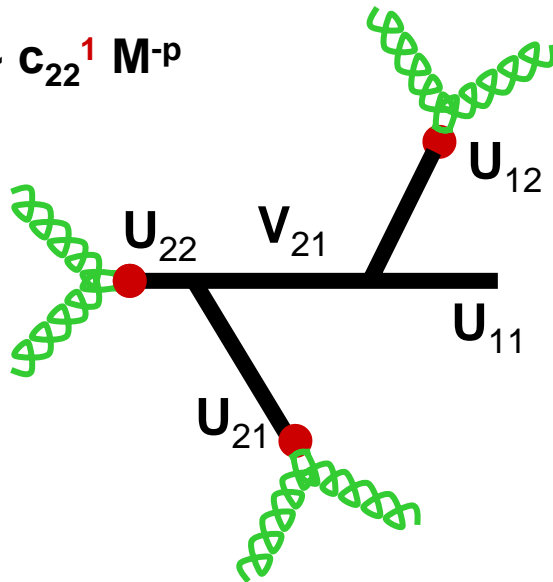
$F \sim 3(N-2)$ (Banks-Zaks weakly coupled)

$$W = m_{JK} Q_J Q_K$$



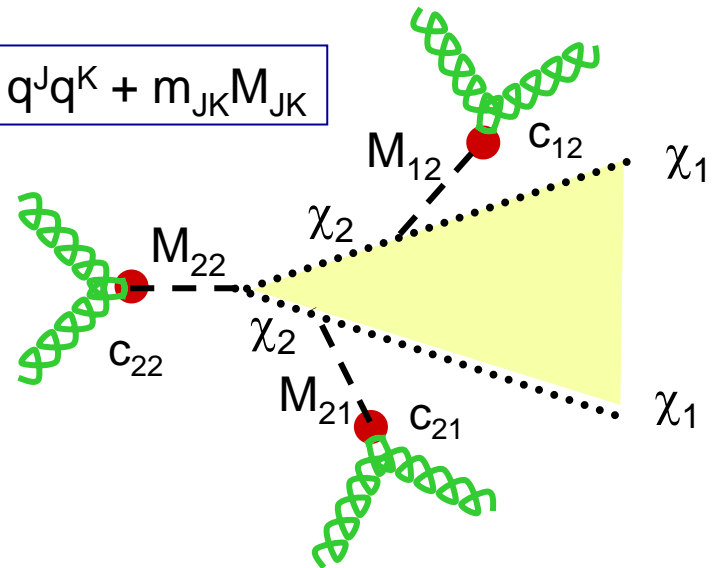
$g g \rightarrow 4 \text{ jets} + \text{MET}$

$$\sim c_{22}^1 M^{-p}$$



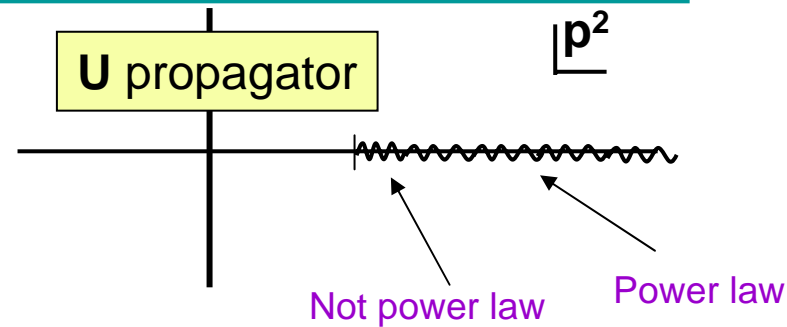
$F \sim \frac{3}{2}(N-2)$ (Seiberg dual weakly coupled)

$$W_{\text{dual}} = M_{JK} q^J q^K + m_{JK} M_{JK}$$



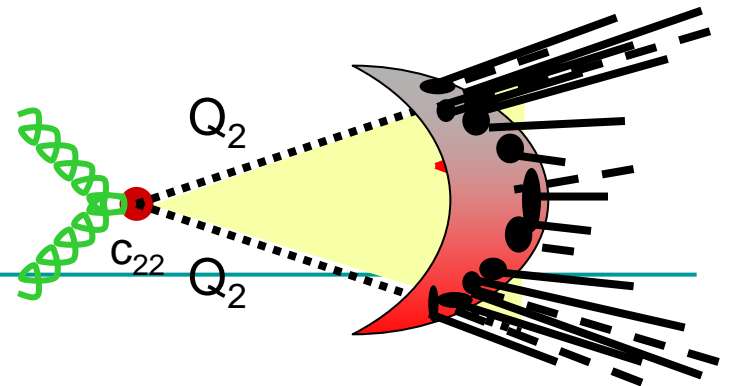
Masses m_{JK} for the Q_J

- Add term $m_{JK}^2 Q_J Q_K^* = m_{JK}^2 \mathbf{V}_{JK}$
- Suppose m_{JK}^2 , c_{JK} simultaneously diagonal
- Mass terms from SUSY breaking \rightarrow too few fermions for new BZ fixed point
- This causes immediate confinement
- v-hadrons could include massless states (v-pions)
- or v-pions could be massive \rightarrow visible decays as in QCD-like models



$g g \rightarrow$ VISIBLE STATES

Visibility at LHC depends on v-pion masses



With a mass gap, there are always poles in some two-point functions

- Put some energy into the vacuum at $x = 0$
- Wait a while
- The energy will eventually distribute itself widely
- There are no long-range forces
- Therefore the final state can be well approximated by a tensor product of localized states, non-interacting with each other, each of which is a stable eigenstate of the Hamiltonian
- Consider any one of these states
- Some local operators interpolate this state
- These operators have poles in their two point functions