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An Introduction to Particle Dark Matter Lecture 1

3rd José Plínio Baptista School on Cosmology

25-30 September 2016 Pedra Azul, ES, Brazil ✓ PhD Theoretical Particle Physics (2004)

International School for Advanced Studies (SISSA-ISAS), Trieste, Italy

- ✓ Postdoc, FSU and California Institute of Technology (2005-2007) *Theoretical Astrophysics and Particle Physics*
- ✓ Joined UCSC Physics Faculty (Assistant Professor, 2007-2011, Associate Professor, July 2011-2015
 Full Professor, July 2015-)
- ✓ Director of UCSC Physics Graduate Studies (2012-)
- ✓ **SCIPP Deputy Director** for **Theory** (July 2011-)



1. What is the origin of the tiny excess of matter over anti-matter?

2. What is the fundamental particle physics nature of Dark Matter?

The Matter Content of the Universe



Dark Matter (83.3%)
 Free H and He (13.8%)
 Stars (1.7%)
 Neutrinos (1.0%)
 Heavy Elements (0,1%)

Please come introduce yourselves!

[to myself, other Instructors, to each other...]

If you are ever on the **US West Coast** please let me know!

Never **underestimate** the importance of **networking** in science!

Feel free to interrupt me during the lectures! (Rare occurrence, but I do occasionally make mistakes...)



C Addison-Wesley Longman



a new elementary particle

...as such it is of interest to particle physicists!

What this mini-lecture series is not:

♦ Review of "evidences" for dark matter

♦ Review of "models" for dark matter

Review of possible claimed "signals" from dark matter (actually, two exceptions...) What this mini-lecture series will be

- ✓ Gross features of dark matter as a particle
- ✓ Paradigms for dark matter in the early universe
- ✓ Schematics of dark matter searches
- ✓ Selected lessons from old and new particle dark matter models

One thing we do **know well** about dark matter

Global amount of dark matter in the **universe**

Reason: very good handles on total energy density, total matter density, total baryonic matter

CMB data indicate the universe is nearly flat
→ energy density is close to critical...

What is the **critical density**? (very good number to have in mind!)

$$ho_{
m crit} \equiv rac{3H_0^2}{8\pi G_N} \simeq 10^{-29} \ {
m g/cm^3}$$

...since 1 GeV ~10⁻²⁴ g, **10 protons** per **cubic meter** (=tiny!)

Various ways to "weigh" matter versus dark energy (CMB+SN+BAO)

...and ordinary (baryonic) matter versus non-baryonic (BBN, CMB) (see Scott Dodelson's lectures)



Global amount of dark matter in the universe from simple subtractions!

$$ar{
ho}_{
m DM} = \Omega_{
m DM}
ho_{
m crit} \simeq 0.3
ho_{
m crit}.$$





clusters... 10⁵ denser!



...which is one of the **key** reasons why **modified gravity** as an alternative to dark matter does not work!

CMB sky is very **boring** – *T* fluctuations very **small**!

T fluctuations proportional to (baryonic) density fluctuations,

$$ar{\delta}
ho/
ho\,\lesssim\,10^{-ar{4}}$$

Matter **over-densities** in linear regime grow **linearly** with scale factor

But the scale factor since CMB decoupling grew by z_{rec}~1,100

Not enough time (since recombination(for structures to go non-linear!

We need a **species** that has **decoupled** from photons much earlier (**Dark Matter**) so that its density **perturbations** are much **larger** at recombination!

 $(\delta \rho / \rho)_{\rm DM} \gg 10^{-4}$

Dark matter **seeds** timely structure formation!







Things go **badly wrong without DM** for structure formation!



Even with best (covariant) incarnation of modified gravity (TeVeS), structure goes non-linear, but the **power spectrum** of matter density fluctuation is **entirely wrong**...



Don't get fooled by the "Vulcan" versus "Neptune" analogy

[Vulcan: No new planet between Mercury and the Sun, but GR Neptune: New planet]

Modified Gravity [MOND, TeVeS] actually does not work at all !!

Knowledge of the dark matter average **density** is a powerful **model-building** tool

Models that **predict** the "right" **amount** of dark matter get kudos

Dark Matter "cosmogony" well-motivated guideline to model building

prototypical example: dark matter as a *thermal relic*... more on this shortly What else do we know about the **microscopic** nature of dark matter from its **macroscopic** features?

"Dark": ...for the reason above! But detailed constraints on electric charge of dark matter are model-dependent... Milli-charge allowed... Phenomenologically: DM is nearly dissipationless (maybe not entirely though, see dark photons, dark disks...)

Collisionless... really? Let's calculate the relevant constraints!



mean free path λ larger than cluster size, ~ 1 Mpc

cluster **density**: $\rho \sim 1 \text{ GeV/cm}^3$, thus...

 $\lambda = 1/(\sigma (\rho/m)) > 1 \text{ Mpc} \rightarrow \sigma/m < 1 \text{ Mpc} / 1 \text{ GeV/cm}^3$

 $\rightarrow \sigma/m < 1 \, cm^2/g$, or 1 barn/GeV



1 barn/GeV... which is strong interaction-size...

is this **small**?

Also, if cross section is **slightly smaller**, no **visible effect**... if cross section **slightly larger**, **disaster**...

Begs the question: is "collisional" **self-interacting** dark matter a "natural" possibility??

Classical: needs to be confined (gravitationally bound) on scales at least as large as dSph... if de Broglie wavelength is larger, disaster strikes!





little exercise: consider $v \sim 100 \text{ km/s}$, show that $\lambda = h/p$ is

$$\lambda \sim 3 \text{ mm} \left(rac{1 \text{ eV}}{m}
ight)$$

which means that to have $\lambda \ll kpc \sim 3x10^{21}$ cm, m>10⁻²² eV

Much, much **better constraints** if the DM is a fermion – we know that the **phase space** density is bounded (Pauli blocking): **f**= **gh**⁻³

Using observed density and velocity dispersion of dSph, Tremaine-Gunn limit (1979): observed phase space density cannot exceed upper bound! (Liouville theorem) Exercise!

 $\sigma \sim 150 \ {
m km/s}$ $ho \gtrsim 1 \ {
m GeV/cm^3}$

$$m^4 > rac{
ho h^3}{[g(2\pi\sigma^2)^{3/2}]} \sim (25 \text{ eV})^4.$$

Fluid: don't want to disrupt pretty (and old!) clusters of stars



Neat exercise to estimate the energy exchanged by encounters of GC and BH, in the impulse approximation, demand that that energy be smaller than binding energy, get maximal mass for BH

Also constraints on **disk stability** ("heating")

Bottom line: *m* < 10³ solar masses ~ 10⁷⁰ eV

...here's the **name of the game**:

(i) Mass: >90 orders of magnitude for bosons, 70 for fermions

(ii) Interactions: ~dark, self-interacting at most ~ strong interactions

(iii) Abundance



Think left and think right and think low and think high.

Oh the things you can think up, if only you try!

Dr. Seuss

A successful framework for the **origin of species** in the early universe: **thermal decoupling**



A successful framework for the **origin of species** in the early universe: **thermal decoupling**



A successful synergy of *statistical mechanics*, *general relativity*, and of *nuclear and particle physics* making **predictions** testable to exquisite accuracy with *astronomical* observations! Key idea of thermal decoupling: if the reaction keeping a species in equilibrium is faster than the expansion rate of the universe, the reaction is in statistical equilibrium; if it's slower, the species decouples ("freeze-out")

$$\Gamma \ll H(T)$$
 $\Gamma(T_{\text{t.o.}}) \sim H(T_{\text{t.o.}})$

the **reaction rate** (from definition of cross section!)

$$\Gamma = n \cdot \sigma \cdot v$$

(1) borrow equilibrium number densities from stat mech

$$n_{
m rel} \sim T^3 \quad {
m for} \ m \ll T,$$

 $n_{
m non-rel} \sim (mT)^{3/2} \exp\left(-rac{m}{T}
ight) \quad {
m for} \ m \gg T.$

(2) borrow Hubble rate from general relativity (FRW solution to Einstein's eq.)

$$H^2 = \frac{8\pi G_N}{3}\rho.$$

$$H^2=rac{8\pi G_N}{3}
ho.$$

GR+SM: energy density in radiation

$$ho \simeq
ho_{
m rad} = rac{\pi^2}{30} \cdot g \cdot T^4 \longrightarrow H \simeq T^2 / M_P$$

first application: **hot** thermal relic

language definition: hot = relativistic at $T_{f.o}$ cold = v < c = 1. (actually not by much, typically!)

simple application: relic SM neutrinos (cosmo v background)

$$\nu + \bar{\nu} \leftrightarrow f + \bar{f},$$

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$$n(T_{\nu}) \cdot \sigma(T_{\nu}) = H(T_{\nu}) \qquad \qquad \sigma \sim G_F^2 T_{\nu}^2$$

suppose this is a hot relic... $n^{-T_v^3}$

$$T_{\nu}^3 G_F^2 T_{\nu}^2 = T_{\nu}^2 / M_P$$

$$T_{\nu} = (G_F^2 M_P)^{-1/3} \simeq (10^{-10} \times 10^{18})^{-1/3} \text{ GeV} \sim 1 \text{ MeV}$$

happy about two things in particular:

1. hot relic assumption works! $T_{\nu} \gg m_{\nu}$

2. Fermi effective theory OK! $T_{\nu} \ll m_W$

$$T_{\nu} = (G_F^2 M_P)^{-1/3} \simeq (10^{-10} \times 10^{18})^{-1/3} \text{ GeV} \sim 1 \text{ MeV}$$

now, how do we calculate the **relic** thermal **abundance** of this prototypical hot relic?

Introduce **Y=n/s** (number and entropy **density**, *V*=*a*³)

If universe is iso-entropic, $s \times a^3 = S$ is conserved

Y ~ n a³ is thus ~ comoving number density, and (without entropy injection)

$$Y_{
m today} = Y_{
m freeze-out} = Y(T_
u)$$

$$Y_{ ext{freeze-out}} = rac{n(T_
u)}{s(T_
u)} = rac{
ho_
u(T_
u)}{m_
u \cdot s(T_
u)}$$

$$Y_{
m today} = Y_{
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$$n_{\mathrm{today}} = s_{\mathrm{today}} imes Y_{\mathrm{today}} = s_{\mathrm{today}} imes Y_{\mathrm{freeze-out}}$$

$$\rho_{\nu, \text{today}} = m_{\nu} \times Y_{\text{freeze-out}} \times s_{\text{today}}$$

$$\Omega_
u h^2 = rac{
ho_
u}{
ho_{
m crit}} h^2 \simeq rac{m_
u}{91.5~{
m eV}}$$

Cowsik-Mc-Clelland limit

That was **fun**! Let's see if it works for something else...

Try **proton-antiproton** freeze-out:

what's the **relic** matter **abundance** in a baryon-symmetric Universe?

$$\sigma \sim \Lambda_{\rm QCD}^{-2}$$

$$n \sigma = H \rightarrow T^3 \Lambda^{-2} = T^2/M_P \rightarrow T = \Lambda^2/M_P$$

doesn't quite work, we're way **outside** the regime of validity for **hot relics**, since T<<<<<m_p...

Need to work out the case of **cold relics**, which looks nastier by eye

$$n \sim (m_{\chi}T)^{3/2} \exp\left(-\frac{m_{\chi}}{T}\right)$$

Here's the trick: **freeze-out** condition gives

$$n_{
m f.o.} \sim rac{T_{
m f.o.}^2}{M_P \cdot \sigma}$$

now define
$$m_{\chi}/T \equiv x$$
 (cold relic: x>>1)

Freeze-out condition (x) now reads

$$rac{m_\chi^3}{x^{3/2}}e^{-x}=rac{m_\chi^2}{x^2\cdot M_P\cdot\sigma}$$

.so we gotta solve
$$\sqrt{x} \cdot e^{-x} = rac{1}{m_\chi \cdot M_P \cdot \sigma}$$

. .

$$\sqrt{x} \cdot e^{-x} = rac{1}{m_\chi \cdot M_P \cdot \sigma},$$





$$\sqrt{x} \cdot e^{-x} = rac{1}{m_\chi \cdot M_P \cdot \sigma},$$

$$\sigma \sim G_F^2 m_\chi^2$$

Take e.g. a "weakly interacting massive particle"

 $m_\chi \sim 10^2$ GeV.

$$\sqrt{x} \cdot e^{-x} = \frac{1}{m_{\chi} \cdot M_P \cdot \sigma} \sim \frac{1}{10^2 \cdot 10^{18} \cdot 10^{-6}} \sim 10^{-14}.$$

thus $x = m_{\chi} / T \sim 35$

Off to calculating the **thermal relic density**

$$\Omega_{\chi} = rac{m_{\chi} \cdot n_{\chi}(T=T_0)}{
ho_c} = rac{m_{\chi} \ T_0^3}{
ho_c} rac{n_0}{T_0^3}$$

iso-entropic universe $aT \sim \text{const}$ $\frac{n_0}{T_0^3} \simeq \frac{n_{\text{f.o.}}}{T_{\text{f.o.}}^3}$

$$\begin{split} \Omega_{\chi} &= \frac{m_{\chi} \ T_0^3}{\rho_c} \frac{n_{\rm f.o.}}{T_{\rm f.o.}^3} = \frac{T_0^3}{\rho_c} x_{\rm f.o.} \left(\frac{n_{\rm f.o.}}{T_{\rm f.o.}^2}\right) = \left(\frac{T_0^3}{\rho_c \ M_P}\right) \frac{x_{\rm f.o.}}{\sigma} \\ & \left(\frac{\Omega_{\chi}}{0.2}\right) \simeq \frac{x_{\rm f.o.}}{20} \left(\frac{10^{-8} \ {\rm GeV}^{-2}}{\sigma}\right) \end{split}$$